

Clustering structure effect on Hanbury-Brown–Twiss correlation in $^{12}\text{C} + ^{197}\text{Au}$ collisions at $\sqrt{s_{NN}} = 200$ GeV

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Outline

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2 AMPT

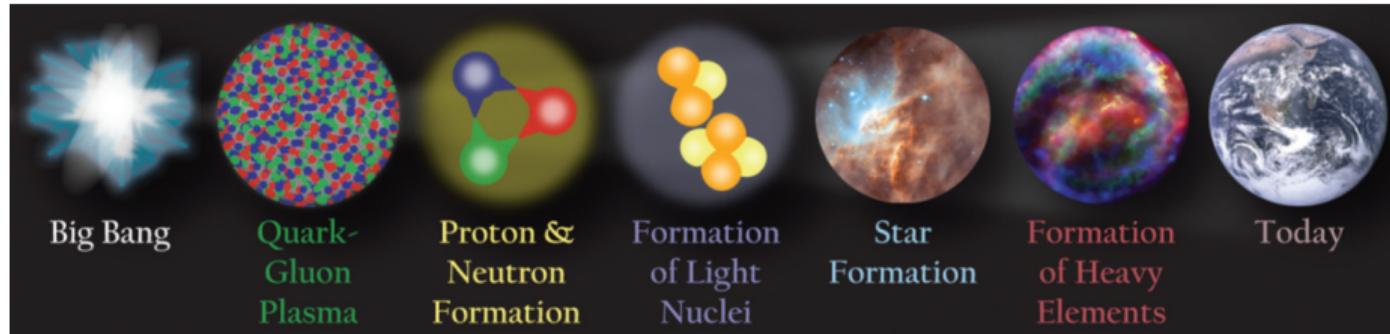
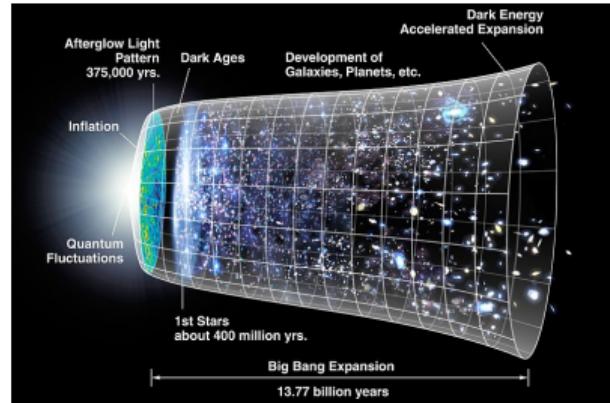
3 HBT

4 Results

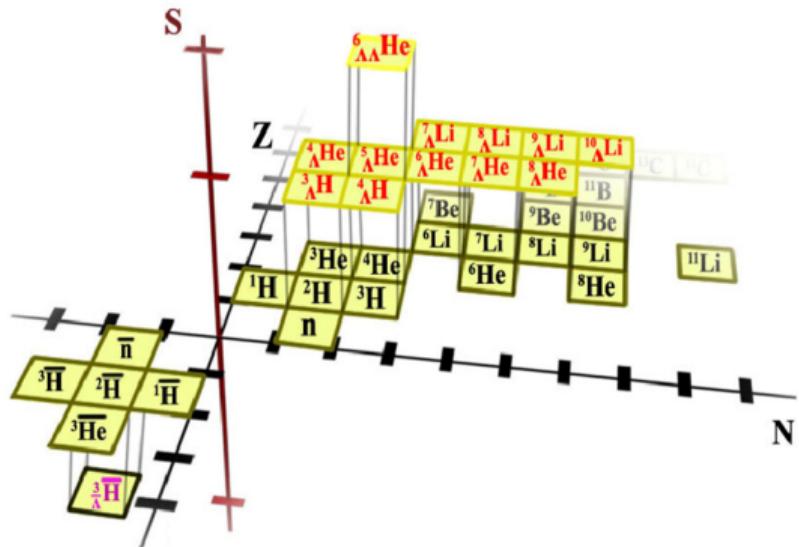
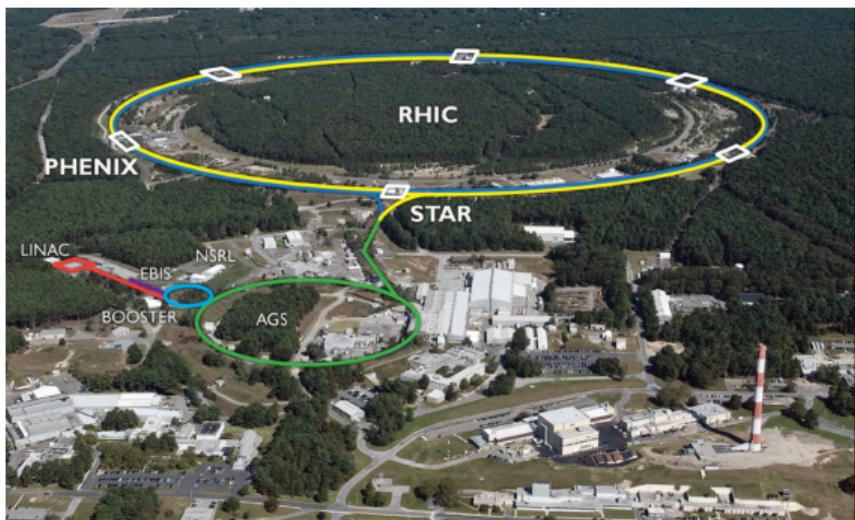
5 Summary

Background

- Quark-Gluon Plasma
- Light Nuclei

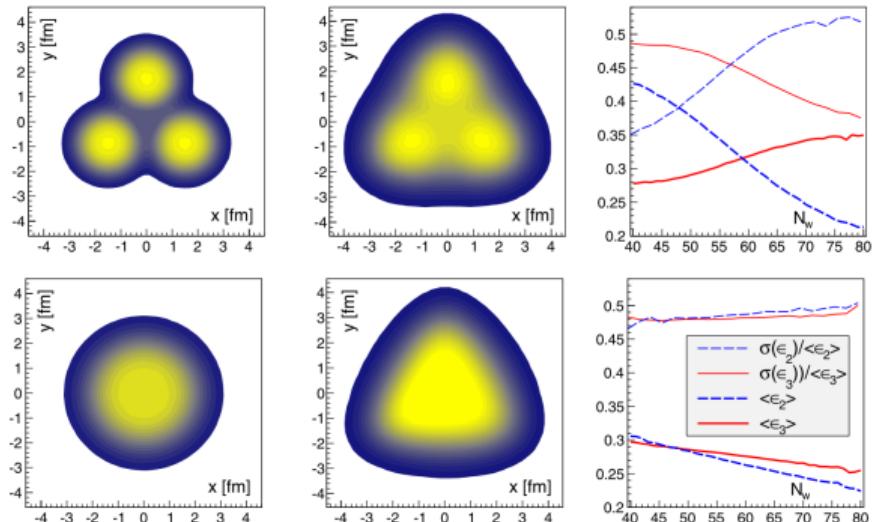


Background

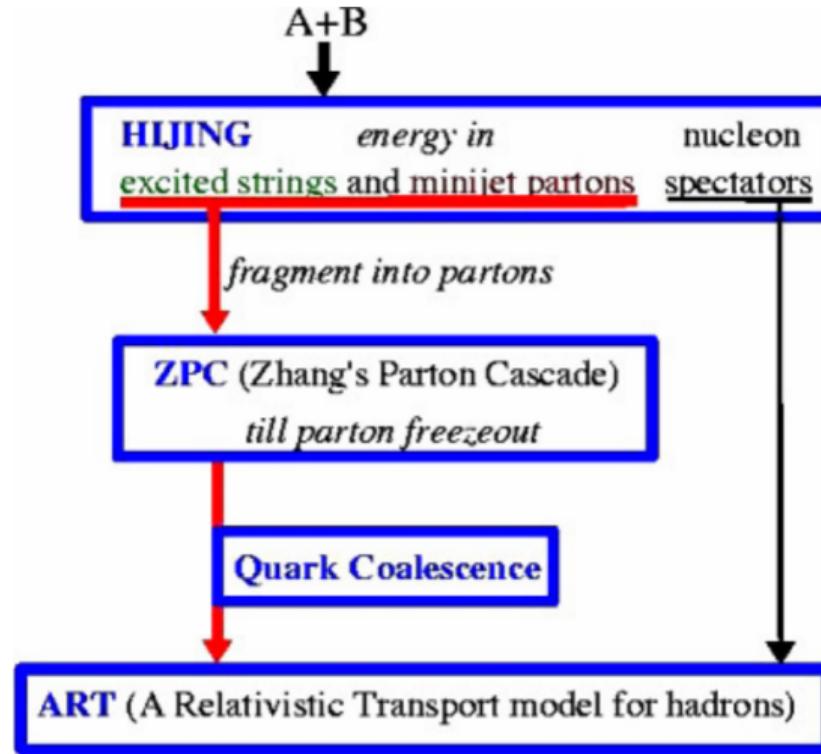


Background

- In 2014, Broniowski, Arriola et al. proposed that through relativistic heavy-ion collision, collective flow can be the signature of α clustering in light nuclei in their ground state

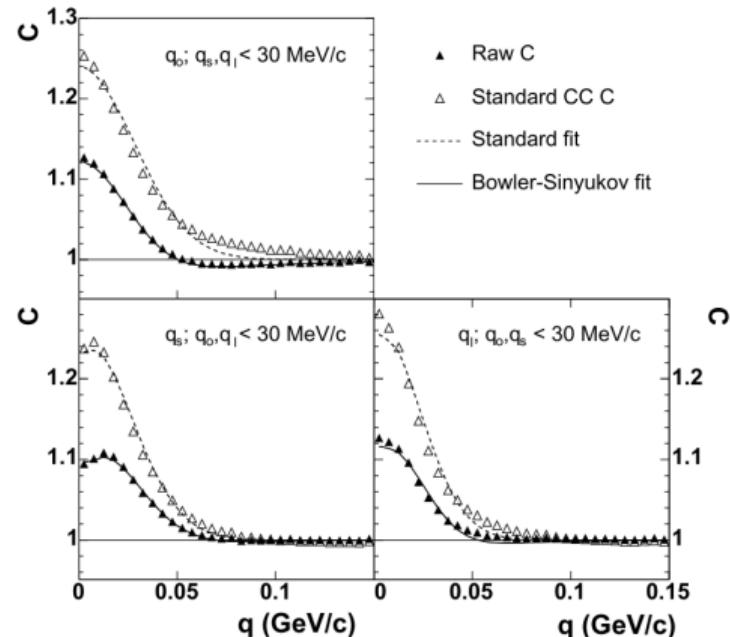
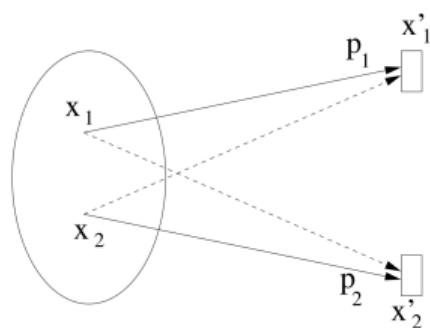


PhysRevLett.112.112501



HBT: History and basic picture

- 1956 Hanbury Brown and Twiss
- 1960 Goldhaber-Goldhaber-Lee-Pais effect



STAR: PhysRevC.71.044906

HBT radii

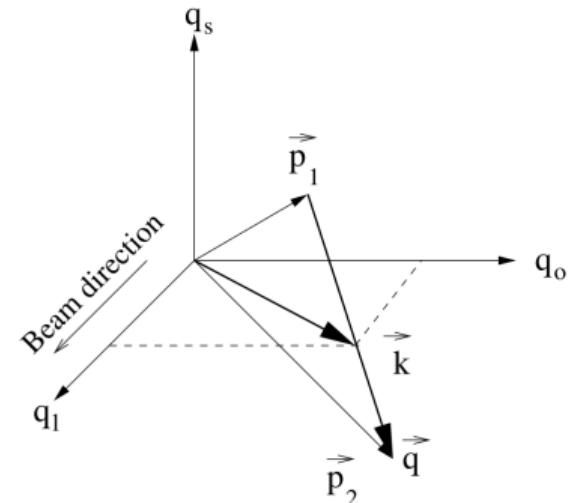
$$C(\vec{q}, \vec{K}) = \frac{\int d^4x_1 d^4x_2 S(x_1, p_1) S(x_2, p_2) |\phi(\vec{q}', \vec{r}')|^2}{\int d^4x_1 S(x_1, p_1) \int d^4x_2 S(x_2, p_2)} \quad (1)$$

$$C(\vec{q}, \vec{K}) = 1 \pm \left| \frac{\int d^4x e^{i\vec{q} \cdot (\vec{x} - \vec{\beta}t)} S(x, K)}{\int d^4x S(x, K)} \right|^2 \quad (2)$$

$$C(\vec{q}, \vec{K}) = 1 \pm e^{-\sum_{i,j=o,s,l} R_{ij}^2(\vec{K}) q_i q_j} \quad (3)$$

$$\begin{aligned} R_s^2(K_T, \Phi, Y) = & \langle (y \cos \Phi - x \sin \Phi)^2 \rangle \\ & - \langle y \cos \Phi - x \sin \Phi \rangle^2 \end{aligned} \quad (4)$$

$$\begin{aligned} R_o^2(K_T, \Phi, Y) = & \langle (x \cos \Phi + y \sin \Phi - \beta_\perp t)^2 \rangle \\ & - \langle x \cos \Phi + y \sin \Phi - \beta_\perp t \rangle^2 \end{aligned} \quad (5)$$



Coulomb Interaction

From Schrödinger Equation with Coulomb potential,

$$\phi_{\text{coulomb}}(\vec{q}, \vec{r}) = \Gamma(1 + i\eta) e^{-\frac{1}{2}\pi\eta} e^{\frac{i}{2}\vec{q} \cdot \vec{r}} F(-i\eta; 1; iz_-) \quad (6)$$

where $\eta = \frac{Z_1 Z_2 \alpha c}{v_{rel}} = \frac{2 Z_1 Z_2 \alpha \mu}{q} = \frac{2}{a_0 q}$, $G(\eta) \equiv e^{-\pi\eta} |\Gamma(1 + i\eta)|^2 = 2\pi\eta/(e^{2\pi\eta} - 1)$ and $z_\pm = \frac{1}{2}(qr \pm \vec{q} \cdot \vec{r})$.

- (anti)symmetrization for identical particles
- Coulomb weight is only valid for $qr/\hbar < 1$

Gaussian Fitting Function

- Quantum Statistics:

$$C(\vec{q}, \vec{K}) = 1 \pm \lambda(\vec{K}) e^{-\sum_{i,j=o,s,l} R_{ij}^2(\vec{K}) q_i q_j} \quad (7)$$

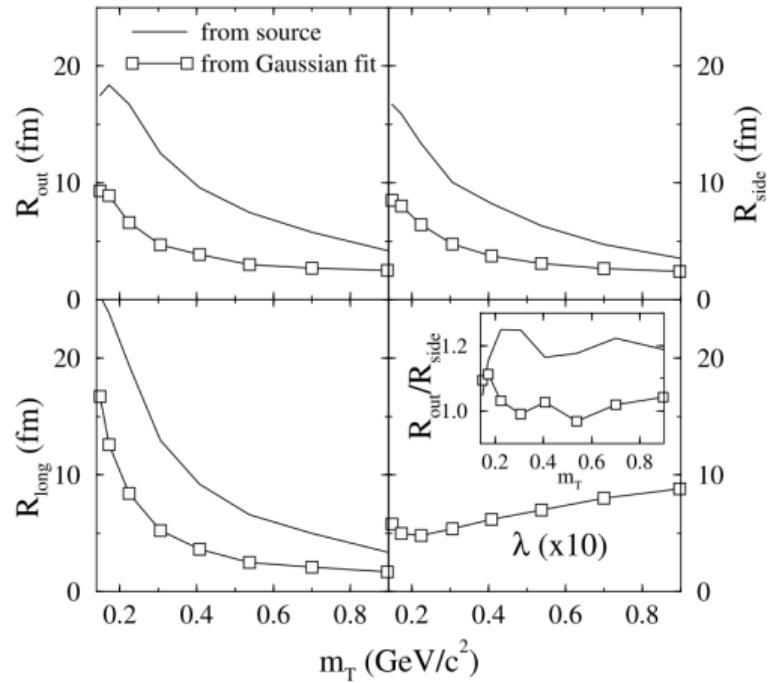
$$\Rightarrow C(q_i) = 1 \pm \lambda e^{-R_i^2 q_i^2} \quad (i = o, s, l, inv) \quad (8)$$

- Coulomb Interaction: (Bowler-Sinyukov procedure)

$$C(\vec{q}, \vec{K}) = (1 - \lambda) + \lambda K_{coul}(q_{inv})(1 \pm e^{-\sum_{i,j=o,s,l} R_{ij}^2(\vec{K}) q_i q_j}) \quad (9)$$

$$\Rightarrow C(q_i) = (1 - \lambda) + \lambda K_{coul}(q_{inv})(1 \pm e^{-R_i^2 q_i^2}) \quad (i = o, s, l, inv) \quad (10)$$

Source radii vs. Gaussian fit



PhysRevLett.89.152301

Q moments and Kurtosis

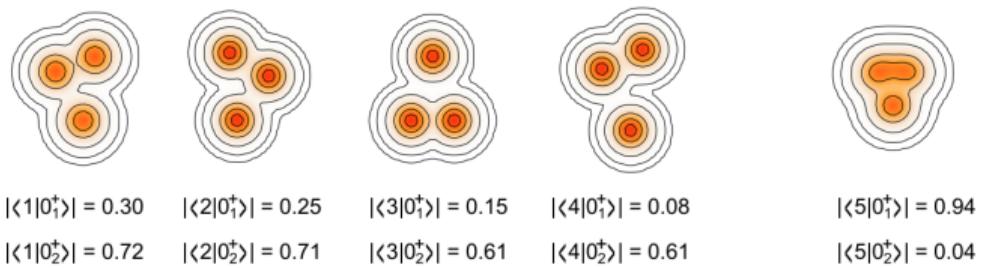
$$\langle\langle q_i^2 \rangle\rangle(\vec{K}) = \frac{\int dq_i q_i^2 [C(\vec{K}, q_i) - 1]}{\int dq_i [C(\vec{K}, q_i) - 1]} \quad (11)$$

$$R_i^2(\vec{K}) = \frac{1}{2 \langle\langle q_i^2 \rangle\rangle} \quad (12)$$

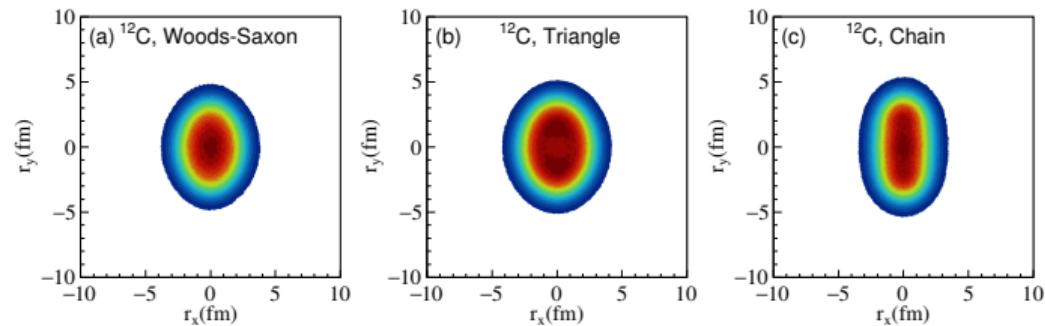
$$\Delta_i(\vec{K}) = \frac{\langle\langle q_i^4 \rangle\rangle(\vec{K})}{3 \langle\langle q_i^2 \rangle\rangle^2(\vec{K})} - 1 \quad i = o, s, l \quad (13)$$

Results: Participant distributions

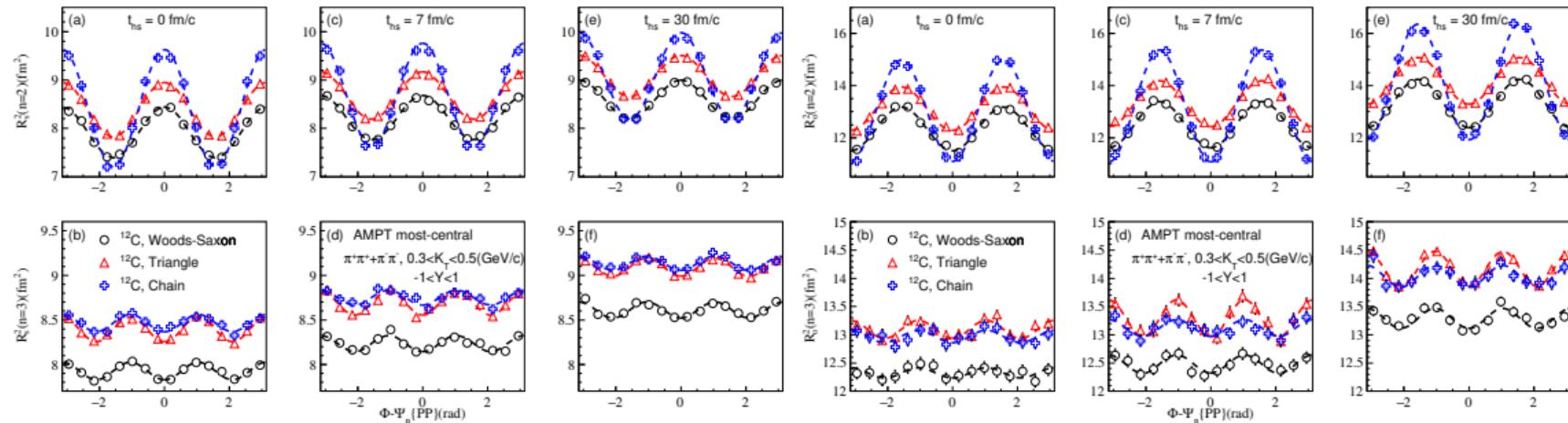
- Chain: Time-dependent Hartree-Fock
- Triangle: Fermionic molecular dynamics, Antisymmetrized molecular dynamics
- *ab initio*



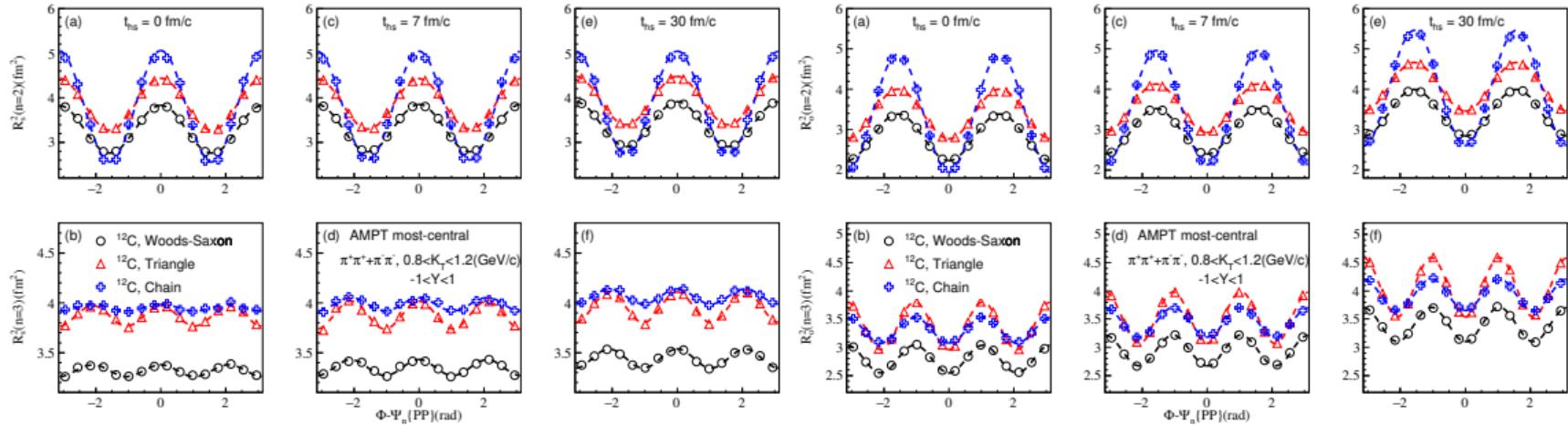
PhysRevLett.98.032501



Azimuthal angle dependence of the HBT radii



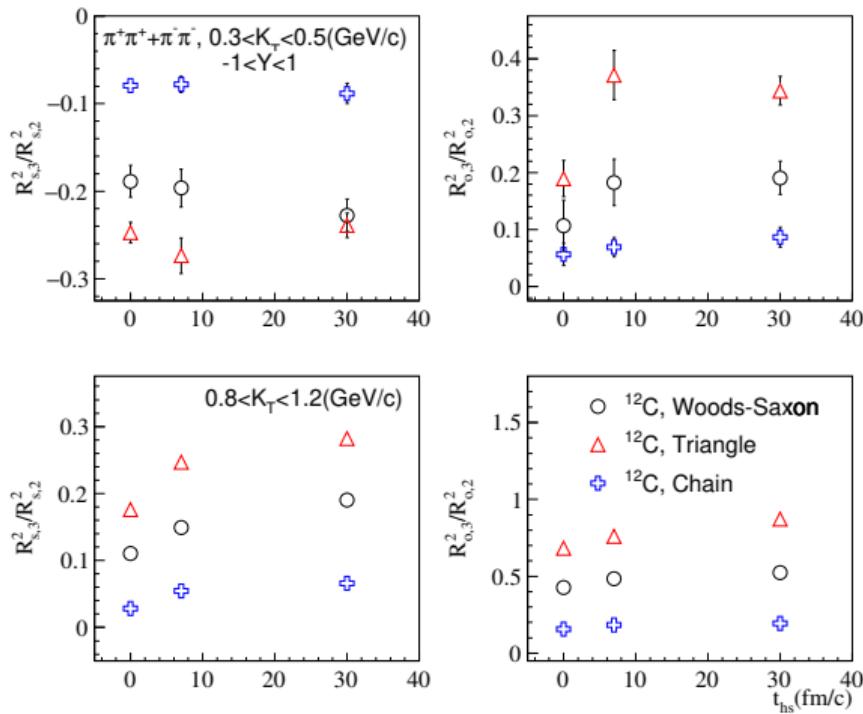
Azimuthal angle dependence of the HBT radii



$$R_s^2(\Phi - \Psi_n) = R_{s,0}^2 + 2R_{s,n}^2 \cos[n(\Phi - \Psi_n)],$$

$$R_o^2(\Phi - \Psi_n) = R_{o,0}^2 + 2R_{o,n}^2 \cos[n(\Phi - \Psi_n)], n = 2, 3 \quad (14)$$

$$R_{s(o),3}^2/R_{s(o),2}^2$$



Summary

- HBT method is useful for studying the geometry and dynamics of fireball
- $R_{s(o),3}^2/R_{s(o),2}^2$ is an effective probe to study the clustering structure of light nuclei
- System scan may provide the platform to further study the bulk properties and exotic structures in light nuclei