

Hengne Li South China Normal University

Experimental Methods and Recent Progresses in Modern Physics





Lecture 2, Hengne Li, SCNU, v1, 2019/03/23



- What it is: a numerical technique for calculating probabilities and related quantities using sequences of random numbers.
- The usal steps:
 - (1) Generate sequence r_1, r_2, \ldots, r_m uniform in [0, 1].
 - (2) Use this to produce another sequence x_1, x_2, \ldots, x_n distributed according to some pdf f(x) in which we are interested (x can be a vector)
- MC generated values = "simulated data", "simulation" ==> use for testing statistical procedures

The Monte Carlo method



 $\int_a^b f(x)\,dx$. Hengne Li, SCNU, 2019/07/17

Random number generators

- Goal: generate uniformly distributed values in [0, 1].
- Tools/algorithms to generate random numbers are called "random number generator"
- There are many algorithms to do that, such as multiplicative linear congruential generator (MLCG), Mersenne twister algorithm, etc.,
- They are implemented, and we can use directly:
 - C++ standard library(stdlib): <u>rand()</u> function can generate random numbers.
 - ROOT: TRandom, TRandom2, TRandom3, where TRandom3 gives much better results [F. James, Comp. Phys. Comm. 60 (1990) 111]

The transformation method

- finding a suitable transformation x(r)
- Remind: previously, we learned given pdf f(x) and function a(x), ask the pdf of a: g(a).
- Here, is given pdf f(x) and g(r), ask the function x(r):

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• Given r_1, r_2, \ldots, r_m uniform in [0, 1], find x_1, x_2, \ldots, x_n that follow f(x) by



F(x) = rHengne Li, SCNU, 2019/07/17

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Example of the transformation method • Exponential pdf: $f(xf(x;\xi) \stackrel{1}{=} \frac{1}{\xi}e^{-x/\xi} \quad (x \ge 0)$ $f(x;\xi) = \frac{1}{\xi}e^{-x/\xi} \quad (x \ge 0)$

Set
$$\int_0^x \frac{1}{\xi} e^{-x'/\xi} \, dx' =$$

$$\rightarrow x(r) = -\xi \ln(1-r)$$

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0

 One has to sol²⁰⁰ to use this met|150 100 arbitrary f(x) 50



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= r and $\int_{0}^{x} \frac{1}{\xi} e^{-x'/\xi} dx' = r$

$$\left(x(r)\right) = -\xi \ln(1-r) \operatorname{WOI}(x(r)) = -\xi \ln r$$



The acceptance-rejection method

- (1) Generate a random number x, uniform in $[x_{\min}, x_{\max}]$, i.e. $x = x_{\min} + r_1 (x_{\max} x_{\min})$, r₁ is uniform in [0, 1].
- (2) Generate a 2nd independent random number u uniformly distributed between 0 and f_{max}, i.e. u=r₂f_{max}
- (3) If u<f(x), then accept x, if not, reject x and repeat
- Practical for f(x) with complicated shapes.



Example with acceptance-rejection method

$$f(x) = \frac{3}{8} (1 \frac{f(x)}{x}) = \frac{3}{8} (1 + x)$$
$$(-1 \le x \le 1)$$

- Generate random number x in [-⁻
- Generate another random numb in [0, 3/4]
- If u <= f(x), take the x, and fill it ir the histogram. If not, drop the x.



Improving efficiency of the acceptance-rejection method

- slow.
- \bullet
 - (1) Generate random variable x according to C h(x) by other method, such as "transformation" method".
 - (2) Generate another random number u uniformly distributted between 0 and C h(x)
 - (3) If u < f(x), take accept x.

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• The fraction of accepted points is equal to the fraction of the box's area under the curve. • For very peaked distributions, the fraction maybe very low, thus the algorithm may be

Improve by enclosing the pdf f(x) in a curve C h(x) that conforms to f(x) more closely, where h(x) is a pdf from which we can generate random values and C is a constant:



Moving to the practical programing training.

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Backup slides, Probability Density Functions.

Probability functions

- We will now introduce a short list of popular probability functions and pdfs.
- For each functions, we show expectation value, variance, a plot and discuss of some properties and applications.
- See also chapter on probability from pdg.lbl.gov

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- Consider N independent experiments (Bernoulli trails):
 - outcome of each is "success" or "failure"
 - probability of "success" on any given trail is p
- Define discrete random variable n = number of successes ($0 \le n \le N$)
- Probability of a specific outcome (in order), e.g. "s s f s f" is: pp(1-p)p(1-p)
- for each permutation.

Binomial distribution

$$) = p^{n}(1-p)^{N-n}$$

• But order is not important; there are N!/[n!(N-n)!] ways (permutations) to get n successes in N trails, total probability for n is sum of probabilities

• The binomial distribution is therefore

$$f(n; N, p) = \frac{N}{n!(N-n!)}$$
random variable parameters

• For the expectation value and variance, we find

$$E[n] = \sum_{n=0}^{N} nf(n; N, p) = Np$$
$$V[n] = E[n^{2}] - (E[n])^{2} = l$$

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Binomial distribution

 $\frac{N!}{(l-n)!}p^n(1-p)^{N-n}$

Np(1-p)

Binomial distribution for different parameters values



 \bullet

Binomial distribution

Example: observe N decays of W bosons, in which n is the number of $W \to \mu v$ decays (number of "success"), and p is the branching ratio (probability of "success").

Multinomia

- Like binomial, but now m outcome $\overrightarrow{p} = (p_1, ..., p_m), \mathbf{w}$
- For N trails, we want the probabili

• This is the multinomial distribution $f(\overrightarrow{n}; N, \overrightarrow{p}) = \frac{N!}{n_1! n_2!}$

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$$\vec{p} = (p_1, \dots, p_m), \quad \text{with } \sum_{i=1}^{n} p_i = 1$$

with $\sum_{i=1}^{m} p_i = 1$
ity to obtain: n_1 of outcome 1,
 n_2 of outcome 2,
 \vdots
 n_m of outcome m.
the for $\vec{n} = (n_1, \dots, n_m)$
 $\frac{n_1!}{\dots n_m!} p_1^{n_1} p_2^{n_2} \cdots p_m^{n_m!} \vec{p} = \frac{N!}{n_1! n_2! \cdots n_m!} p_1$





Multinomial distribution

- Now consider outcome i as "success", all others are "failure". $E[n_i] = Np_i, \quad V[n_i] = Np_i(1-p_i)$ for all *i*
 - Note, i (success) and NOT-i (failure), actually form a "binomial" structure
- One can also find the covariance to be $V_{ij} = N p_i \left(\delta_{ij} p_j \right)$
- Example: $\vec{n} = (n_1, ..., n_m)$ represents a histogram with m bins, N total entries, all entries independent.

Poisson distribution

 Consider binomial n in the limit $N \to \infty, \quad p \to 0,$

- n follows the Poisson distribution: $f(n;\nu) = \frac{\nu^n}{n!} e^{-\nu} \quad (n \ge 0)_{E[n] = \nu},$ $E[n] = \nu, \quad V[n] = \nu$
- Example: number of scattering events n with cross section σ found for a fixed integrated luminosity, with $\nu = \sigma | Ldt$

 $N \to \infty, \qquad p \to 0, \qquad E[n] = Np \to \nu.$

$$E[n] = Np \to \iota$$

$$e^{-\nu}$$
 $(n \ge 0)$

$$V[n] = \nu .$$



Uniform distribution

• Consider a continuous random vari $f(x; \alpha, \beta) = \begin{cases} \frac{1}{\beta - \alpha} < \alpha \leq x \leq \beta \\ 0 & \text{otherwise} \end{cases}$ Uniform pdf is

$$f(x; \alpha, \beta) = \begin{cases} \frac{1}{\beta - \alpha} & \alpha \le x \le \beta \end{cases} \begin{cases} \widehat{\beta} & E \\ 0 & \text{otherwise} \end{cases} \begin{cases} \widehat{\beta} & V \\ V & V \end{cases}$$

• $NE[x] = \frac{1}{2}(\alpha + \beta)d\sigma[m] variable + distribution F(x), y = F(x) is unifor$ $<math>V[x] = \frac{1}{12}(\beta - \alpha)^2 [x] + E_{\gamma 1} + Unifc$

$$E_{\min} = \frac{1}{2} E_{\pi} (1 - \beta), \quad E_{\max} =$$

 $E_{\text{min}} = \frac{1}{-}E_{\pi}(1-\beta)$ $E_{\text{max}} = \frac{1}{-}E_{\pi}(1+\beta)$



 $\beta = v/c$

Exponential distribution

- $\begin{array}{c} \textbf{0.4} \\ \textbf{E[x]} \\ \textbf{0.2} \end{array}$ $E[x] = E[x] = \xi$ $V[x] = \hat{V[x]} = \xi^2$ V[x]
- Example: proper decay time t of an unstable particle $f(t; \tau) = \frac{1}{\tau} e^{-t/\tau}$ (τ = mean lifetime) τ

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$$\int f(t) = f(t)$$

$$f(t - t_0 | t \ge t_0) = f(t) \quad f(t - t_0 | t \ge t_0) = f(t)$$

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$$f(x;\mu,\sigma) = \frac{1}{\sqrt{2\pi\sigma}} e^{-(x-\mu)^2/2\sigma^2}$$

 $E[x] = \mu$ (N.B. often μ , σ^2 denote mean, variance of any $V[x] = \sigma^2$ r.v., not only Gaussian.)

• Special case:

 $V[x] = \sigma^{--} \mathbf{S}$ $\mu = 0, \, \sigma^2 = 1$

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Gaussian distribution

• The Gaussian (normal) pdf for a continuous random variable x is defined by





Gaussian pdf and the Central Limit Theorem

- Central Limit Theorem:
- arbitrary pdfs, consider the sum

• In the limit $n \to \infty$, y is a Gaussian random variable with

$$E[y] = \sum_{i=1}^{n} \mu_i \qquad V[y] = \sum_{i=1}^{n} \sigma_i^2 \quad V[y] = \sum_{i=1}^{n} \sigma_i^2$$

measured values can be treated as Gaussian random variables.

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 The Gaussian pdf is very useful because almost any random variable that is a sum of a large number of small contributions follows it. This follows from the

• For n independent random variables x_i with finite variances σ_i^2 , otherwise $=\sum_{i=1}^{n} x_{i}$ i=1

Measurement errors are often the sum of many contributions, so frequently

Central Limit Theorem (CLT)

- The CLT can be proved using characteristic functions (Fourier transforms).
- For finite n, the theorem is approximately valid to the extent that the fluctuation of the sum is not dominated by one (or few) terms.
 => beware of measurement errors with non-Gaussian tails.
- Good example: velocity components v_x of air molecules.
- OK example: total deflection due to multiple Coulomb scattering. (Rare large angle deflections give non-Gaussian tail.)
- Bad example: energy loss of charged particle traversing thin gas layer.

Multivariate Gaussian $d_{\vec{x}} = (x_1, \dots, x_n)^{t}$ ion

- Multivariat $\cap \bigcap_{i=1}^{\infty} f(\vec{x}; \vec{\mu}, V) = \frac{1}{(2\pi)^2}$ $f(\vec{x};\vec{\mu},V) = \frac{1}{(2\pi)^{n/2}} V^{1/2} V^{1/2}$ • $\vec{x}, \vec{\mu} \in \vec{x}, \vec{\mu}$ double of the set of the s $E E[x_i] = \mu_i, \quad \text{COV}[x_i, x_j]$

• For n=2 this is
$$f(x_1, x_2; ; \mu_1, \mu_2, \sigma_1, \sigma_2, \rho) = \frac{1}{2\pi\sigma_1\sigma_2\sqrt{1-\rho^2}}$$
$$f(x_1, x_2; ; \mu + \mu_2, \sigma_1, \sigma_2, \rho) = \frac{1}{2\pi\sigma_1\sigma_2\sqrt{1-\rho^2}}$$
$$(x_1, x_2; ; \mu + \mu_2, \sigma_1, \sigma_2, \rho) = \frac{1}{2\pi\sigma_1\sigma_2\sqrt{1-\rho^2}}$$
$$(x_1, x_2; ; \mu + \mu_2, \sigma_1, \sigma_2, \rho) = \frac{1}{2\pi\sigma_1\sigma_2\sqrt{1-\rho^2}}$$
$$(x_1, x_2; ; \mu + \mu_2, \sigma_1, \sigma_2, \rho) = \frac{1}{2\pi\sigma_1\sigma_2\sqrt{1-\rho^2}}$$
$$(x_1, x_2; ; \mu + \mu_2, \sigma_1, \sigma_2, \rho) = \frac{1}{2\pi\sigma_1\sigma_2\sqrt{1-\rho^2}}$$
$$(x_1, x_2; ; \mu + \mu_2, \sigma_1, \sigma_2, \rho) = \frac{1}{2\pi\sigma_1\sigma_2\sqrt{1-\rho^2}}$$
$$(x_1, x_2; ; \mu + \mu_2, \sigma_1, \sigma_2, \rho) = \frac{1}{2\pi\sigma_1\sigma_2\sqrt{1-\rho^2}}$$
$$(x_1, x_2; ; \mu + \mu_2, \sigma_1, \sigma_2, \rho) = \frac{1}{2\pi\sigma_1\sigma_2\sqrt{1-\rho^2}}$$
$$(x_1, x_2; ; \mu + \mu_2, \sigma_1, \sigma_2, \rho) = \frac{1}{2\pi\sigma_1\sigma_2\sqrt{1-\rho^2}}$$
$$(x_1, x_2; ; \mu + \mu_2, \sigma_1, \sigma_2, \rho) = \frac{1}{2\pi\sigma_1\sigma_2\sqrt{1-\rho^2}}$$
$$(x_1, x_2; ; \mu + \mu_2, \sigma_1, \sigma_2, \rho) = \frac{1}{2\pi\sigma_1\sigma_2\sqrt{1-\rho^2}}$$
$$(x_1, x_2; ; \mu + \mu_2, \sigma_1, \sigma_2, \rho) = \frac{1}{2\pi\sigma_1\sigma_2\sqrt{1-\rho^2}}$$
$$(x_1, x_2; ; \mu + \mu_2, \sigma_1, \sigma_2, \rho) = \frac{1}{2\pi\sigma_1\sigma_2\sqrt{1-\rho^2}}$$
$$(x_1, x_2; ; \mu + \mu_2, \sigma_1, \sigma_2, \rho) = \frac{1}{2\pi\sigma_1\sigma_2\sqrt{1-\rho^2}}$$
$$(x_1, x_2; ; \mu + \mu_2, \sigma_1, \sigma_2, \rho) = \frac{1}{2\pi\sigma_1\sigma_2\sqrt{1-\rho^2}}$$
$$(x_1, x_2; ; \mu + \mu_2, \sigma_1, \sigma_2, \rho) = \frac{1}{2\pi\sigma_1\sigma_2\sqrt{1-\rho^2}}$$

$$\frac{1}{\pi^{n/2}|V|^{1/2}} \exp\left[-\frac{1}{2}(\vec{x}-\vec{\mu})^T V^{-1}(\vec{x}-\vec{\mu})\right]$$

$$\exp\left[-\frac{1}{2}(\vec{x}-\vec{\mu})^T V^{-1}(\vec{x}-\vec{\mu})\right]$$

$$= V_{ij}$$
.

$$n = 1, 2, \dots = \text{number of 'degrees of}$$
$$E[z] = n^{n-1} \widetilde{V}[z] \stackrel{i}{=} 2n .$$
$$E[z] = n, \quad V[z] = 2n .$$

• For independent Gaussian x_i, $i \stackrel{E[z]}{=} n, \quad V[z] = 2n$.

$$z = z = \sum_{i=1}^{n} \frac{(x_i - \mu_i)^2}{\sigma_i^2}$$
 fol

least squares.

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Cauchy (Breit-Wigner) distribution

- E[x] is not well defined, $V[x] \to \infty^{\circ}$ $x_0 = mode$ (most probable value) Γ = full width at half maximum
- Example: mass of resonance particles, e.g. $Z^0, \rho, K^*, \phi^0, \dots$
- $\Gamma = \text{decay rate}$ (inverse of mean lifetime)





Landau distribution

d, the energy loss Δ follows the Landau pdf:

$$f(\Delta;\beta) = \frac{1}{\xi}\phi(\lambda) ,$$

$$\phi(\lambda) = \frac{1}{\pi} \int_0^\infty \exp(-u \ln u)$$

$$\lambda = \frac{1}{\xi} \left[\Delta - \xi \left(\ln \frac{\xi}{\epsilon'} + 1 - \frac{2\pi N_{\rm A} e^4 z^2 \rho \sum Z}{m_{\rm e} c^2 \sum A} \frac{d}{\beta^2} \right) \right]$$

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• For a charged particle with $\beta = v/c$ traversing a layer of matter of thickness



Landau distribution

Long 'Landau tail' \rightarrow all moments ∞

Mode (most probable value) sensitive to β , \rightarrow particle i.d.

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$$f(x; \alpha, \beta) = \frac{\Gamma(\alpha + \beta)}{\Gamma(\alpha)\Gamma(\beta)} x^{\alpha-1} (1 - x)^{\beta-1} x^{\beta-1} (1 - x)^{\beta-1} x^{\alpha-1} (1 - x)^{\beta-1} x^{\beta-1} (1 - x)^{\beta-1} (1 - x)^{\beta-1} x^{\beta-1} (1 - x)^{\beta-1} (1 - x)^{\beta-1} x^{\beta-1} (1 - x)^{\beta-1} x^{\beta-1} (1 - x)^{\beta-1} (1 - x)^{\beta-1}$$

Often used to respreser 0.5
 of continuous non-zero 00
 random variable betwecon
 finite limits

$c^{\alpha-1}(1-x)^{\beta-1}$ ion



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0.2

0

0

||

$$f(x; \alpha, \beta) = \frac{1}{\Gamma(\alpha)\beta^{\alpha}} x^{\alpha-1} e^{-x/\beta} e^{-x/\beta}$$

[0,∞].

 Also e.g. sum of n exponentia ^{0.1} random variables or time until event in Poisson process ~ ⁰
 Gamma

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$x^{\alpha-1}e^{-x/\beta}$ but ion

 $|\beta|$



$$\mathbf{S}_{f}(x;\nu) = \frac{\Gamma\left(\frac{\nu+1}{2}\right)}{\sqrt{\nu\pi}\Gamma(\nu/2)}}{\sqrt{\nu\pi}\Gamma(\nu/2)}$$

$$f(x;\nu) = \frac{\Gamma\left(\frac{\nu+1}{2}\right)}{\sqrt{\nu\pi}\Gamma(\nu/2)}}{\sqrt{\nu\pi}\Gamma(\nu/2)}\left(1 + \frac{x^{2}}{\nu}\right)^{-\left(\frac{\nu+1}{2}\right)}}{\nu} \int_{-\left(\frac{\nu+1}{2}\right)}^{-\left(\frac{\nu+1}{2}\right)}}{\sqrt{\nu\pi}\Gamma(\nu/2)} \int_{-\left(\frac{\nu+1}{2}\right)}^{-\left(\frac{\nu+1}{2}\right)}}{\nu} \int_{-\left(\frac{\nu+1}{2}\right)}}{\nu} \int_{-\left(\frac{\nu+1}{2}\right)}}{\nu} \int_{-\left(\frac{\nu+1}{2}\right)}^{-\left(\frac{\nu+1}{2}\right)}}{\nu} \int_{-\left(\frac{\nu+1}{2}\right)}}{\nu} \int_{-\left(\frac{$$

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Student's t distribution

- If x ~ Gaussian with $\mu = 0$, $\sigma^2 = 1$, and $z \sim \chi^2$ with n degrees of freedom, then $t = x / (z/n)^{1/2}$ follows Student's t with v = n.
- This arises in problems where one forms the ratio of a sample mean to the sample standard deviation of Gaussian random variables.
- The Student's t provides a bell-shaped pdf with adjustable tails, ranging from those of a Gaussian, which fall off very
- quickly, $(v \rightarrow \infty)$, but in fact already very Gauss-like for v = two dozen), to the very long-tailed Cauchy (v = 1).

Monte Carlo eve

- Simple example: $e^+e^- \rightarrow \mu^+\mu^-$
- Generate $\cos\theta$ and ϕ $f(\cos\theta; A_{\text{FB}}) \propto (1 + \frac{8}{3}A_{\text{FB}}\cos\theta + \cos^2\theta)$,
 - $f(g(\phi) = \frac{1}{2\pi} \quad (0 \le \phi \le 2\pi)$
- Let $g(f(\cos\theta; A_{\text{FB}}) \propto (1 + \frac{8}{3}A_{\text{FB}}\cos\theta + \cos^2\theta),$ $g(\phi) = \frac{1}{2\pi} \quad (0 \le \phi \le 2\pi)$
- e.g. PYTHIA, POWHEG, HERWIG, ISAJET ...
- Output = "events", and for each even their momentum vectors, types, etc..

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Output = "events", and for each event we get a list of generated particles and

Example of a simulated event

\times ~							
Event listing (summary)			A simulated event				
I particle/jet KS KF orig p_x p_y p_z E	m						
1 !p+! 21 2212 0 0.000 0.000 7000.000 7000. 2 lp+! 21 2212 0 0.000 0.000-7000.000 7000.	000 0,938						
Z (p) 21 2212 0 0,000 0,000 1000,000 1000, Z (p) 21 21 1 0.863 -0.323 1739 862 1739	862 0.000			•			
4 !ubar! 21 -2 2 -0.621 -0.163 -777.415 777.	415 0,000						
5 !g! 21 21 3 -2,427 5,486 1487,857 1487. 6 lol 21 21 4 -62,910 63,357 -463,274 471	Χ~					- D ×	
7 !~9! 21 1000021 0 314,363 544,843 498,897 979	397 pi+	1 211	209 0,006	0.398 -308.296	308,297	0,140	
8 !"g! 21 1000021 0 -379,700 -476,000 525,686 980.	398 gamma	1 22	211 0,407	0,087-1695,458 :	1695,458	0,000	
9 !"ch1_1-! 21-1000024 7 130.098 112.247 129.860 263. 10 Isbarl 21 -3 7 259 400 187 468 83 100 330	400 (pi0) 1	1 22	212 0.021	0.122 -103.709	103.709	0.135	
11 lc! 21 4 7 -79.403 242.409 283.026 381	401 (pi0) 1	11 111	212 0,084	-0,068 -94,276	94,276	0,135	
12 !"chi_20! 21 1000023 8 -326,241 -80,971 113,712 385.	402 (pi0) 1	11 111	212 0,267	-0,052 -144,673	144.674	0,135	
13 !b! 21 5 8 -51.841 -294.077 389.853 491.	403 gamma	1 22	215 -1,581	2,473 3,306	4,421	0.000	
14 !bban! 21 -5 8 -0,597 -99,577 21,299 101.	404 gamma 405 pi-	1 22	215 -1,494	2,143 5,051 0 779 4 015	4.015	0.000	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	405 pi+	1 211	216 -0.024	0.293 0.486	0.585	0.140	
$17 \ cbar 21 -4 9 \ 20.839 \ -7.250 \ -5.938 \ 22$	407 K+	1 321	218 4.382	-1.412 -1.799	4,968	0.494	
18 !"chi_10! 21 1000022 12 -136,266 -72,961 53,246 181.	408 pi-	1 -211	218 1,183	-0,894 -0,176	1,500	0,140	
19 !nu_mu! 21 14 12 -78,263 -24,757 21,719 84.	409 (pi0) 1	11 111	218 0,955	-0,459 -0,590	1,221	0,135	
20 !nu_mubar! 21 -14 12 -107.801 16.901 38.226 115.	410 (pi0) 1	11 111	218 2,349	-1,105 -1,181	2,855	0,135	
21 or 21 or	411 (KDarv) 1 412 pi-	1 -011	219 1,441 219 2,222	-0,247 -0,472	2 225	0,498	
21 gamma 1 22 4 $2,656$ $1,557$ $0,125$ 2.22 ("chi 1-) 11-1000024 9 129 643 112 440 129 820 262	412 p1 413 K+	1 321	220 1.380	-0.652 -0.361	1.644	0.494	
23 ("chi 20) 11 1000023 12 -322.330 -80.817 113.191 382	414 (pi0) 1	11 111	220 1.078	-0.265 0.175	1.132	0.135	
24 ~chi_10 1 1000022 15 97,944 77,819 80,917 169.	415 (K_SO) 1	11 310	222 1,841	0,111 0,894	2,109	0,498	
25 "chi_10 1 1000022 18 -136,266 -72,961 53,246 181.	416 K+	1 321	223 0,307	0,107 0,252	0,642	0,494	
26 nu_mu 1 14 19 -78,263 -24,757 21,719 84.	417 pi-	1 -211	223 0,266	0.316 -0.201	0,480	0.140	
27 nu_mubar 1 -14 20 -107,801 16,901 38,226 115.	418 nbar0	1 -2112	226 1,535	1.641 2.078	3,111 1 000	0.940	
28 (Delta++) 11 2224 2 0.222 0.012-2/34.28/ 2/34.	419 (p10) 1 420 pit	1 211	226 0,633 227 0.217	1,046 1,511	1,908	0,135	
<u>a</u>	421 (pi0) 1	11 111	227 1.207	2.336 2.767	3,820	0.135	
	422 n0	1 2112	228 3.475	5.324 5.702	8,592	0.940	
•	423 pi-	1 -211	228 1,856	2,606 2,808	4,259	0,140	
•	424 gamma	1 22	229 -0,012	0,247 0,421	0,489	0.000	
•	425 gamma	1 22	229 0.025	0.034 0.009	0.043	0.000	
	425 p1+	1 211	250 2,718	5,229 5,403	8,703 10.901	0,140	
	427 (p10)	1 -211	230 4,103	1.233 1.945	2 372	0.140 ***	
PYIHIA Monte Carlo	429 (pi0) 1	11 111	231 0.645	1.141 0.922	1.608	0.135	
	430 gamma	1 22	232 -0,383	1,169 1,208	1,724	0,000	
$pp \rightarrow gluino-gluino$	431 gamma	1 22	232 -0,201	0,070 0,060	0,221	0,000	

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Monte Carlo detector simulation

- Takes as input the particle list and momenta from generator.
- Simulates detector response:
 - hadronic showers, production of signals, electronics response, ...
- Output = simulated raw data -> input to reconstruction software: track finding, fitting, shower clustering, etc.
- "efficiency" = [N events found] / [N events generated]
- Software: GEANT4

• multiple Coulomb scattering (generate scattering angle), particle decays (generate lifetime), ionization energy loss (generate Δ), electromagnetic,

 Predict what you should see at "detector level" of events generated by the "event generator". Compare with "real data", e.g. can be used to estimate