

# Precision test of Standard Model and search for New Physics using lattice QCD

Xu Feng (冯旭)



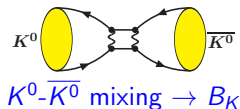
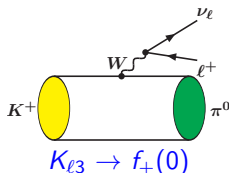
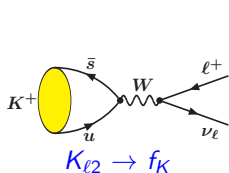
Frontier of QCD: Opportunities and Challenges @ PKU, 2019/11/10

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恭祝赵老师八十华诞  
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# Role of lattice QCD in flavor physics

Lattice QCD is powerful for observables such as



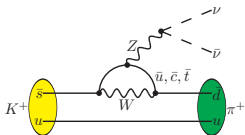
Flavor Lattice Averaging Group (FLAG) average 2019

	$N_f$	FLAG average	Frac. Err.
$f_K/f_\pi$	2 + 1 + 1	1.1932(19)	0.16%
$f_+(0)$	2 + 1 + 1	0.9706(27)	0.28%
$f_{D_s}/f_D$	2 + 1 + 1	1.1783(16)	0.13%
$\hat{B}_K$	2 + 1	0.7625(97)	1.27%

lattice QCD calculations play important role in precision flavor physics

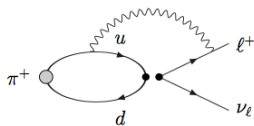
## Search for New Physics in rare processes / high-intensity frontiers

### 2<sup>nd</sup> order electroweak interaction



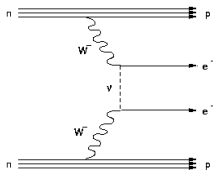
$$\text{Br}[K^+ \rightarrow \pi^+ \nu \bar{\nu}] = 1.73_{-1.05}^{+1.15} \times 10^{-10}$$

### Electromagnetic correction



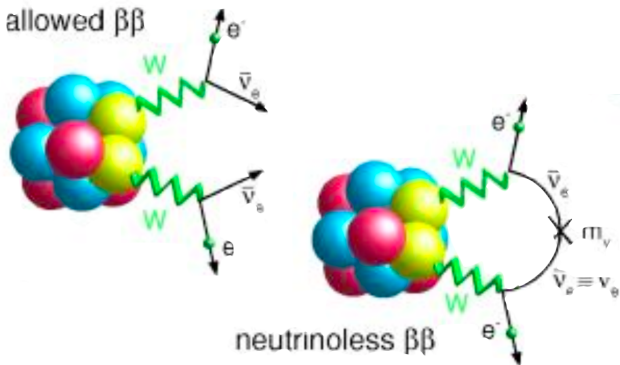
$$f_K/f_\pi = 1.1932(19) \Rightarrow 0.16\% \text{ uncertainty}$$

### Neutrinoless double beta decay



$$T_{1/2}^{0\nu} > 10^{25} \text{ year}$$

# $0\nu 2\beta$ decays



Inspired by a talk given by Prof. HuanZhong Huang

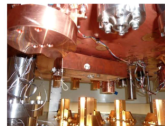
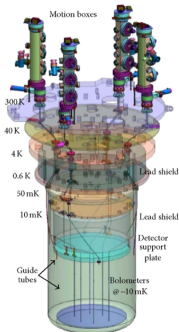
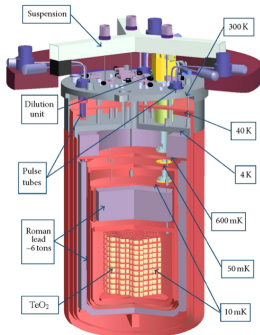


## Majorana Neutrinos &



## Recent CUORE Result on Neutrinoless Double Beta Decay

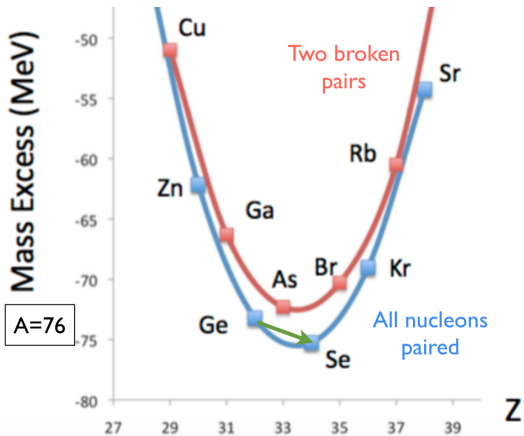
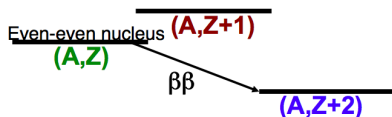
Huan Zhong Huang (黄焕中)



# Double beta decays

## Early in 1935, Goppert-Mayer propose to detect double beta decay

- Nuclear pairing: In some case even-even nucleus is more stable, e.g.  $\text{Ge}^{76}$

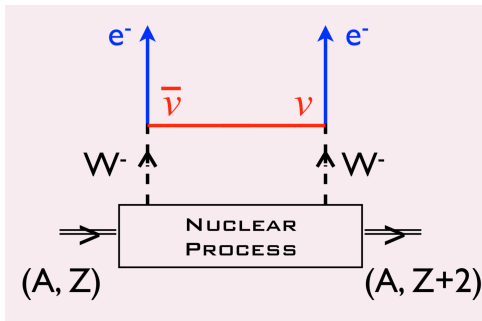


# Majorana neutrinos

## Majorana's proposal in 1937: $\nu = \bar{\nu}$ ?

⇐ This is allowed by symmetry properties of Dirac's theory

- In single beta decay, one cannot distinguish Dirac or Majorana neutrino
- 1939, Furry propose to search for neutrinoless double beta ( $0\nu\beta\beta$ ) decays



- The process violates the lepton number by two units

Question: do we need the lepton number conservation?



# Lepton number conservation

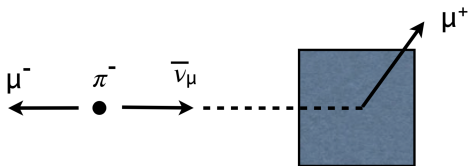
According to phase space factor,  $0\nu\beta\beta$  mode is highly favored over  $2\nu\beta\beta$

$$T_{1/2}^{2\nu 2\beta} \approx 10^{25} \text{ yr}, \quad T_{1/2}^{0\nu 2\beta} \approx 10^{19} \text{ yr}$$

However

- $2\nu\beta\beta$  has been detected in total of 10 nuclei:  $^{48}\text{Ca}$ ,  $^{76}\text{Ge}$ , ...  $^{238}\text{U}$
- No  $0\nu\beta\beta$  detected yet

Also, in neutrino capture,  $\bar{\nu}$  always produce positive charged lepton

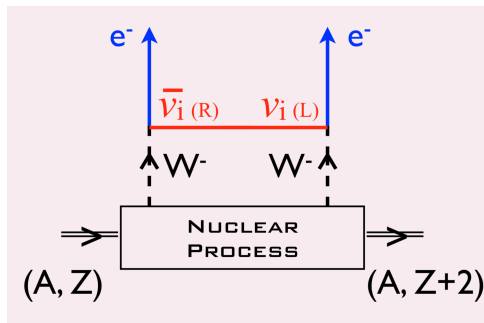


Consequence: Introduce lepton number conservation to explain experiments

# Maximal Parity Violation

1956, Lee & Yang discover parity violation in weak decays  
[Nobel prize 1957]

- Neutrino is left-handed, while anti-neutrino is right-handed
- Helicity exactly forbids the second vertex in  $0\nu\beta\beta$  already
  - ▶ Lepton number conservation is no longer needed

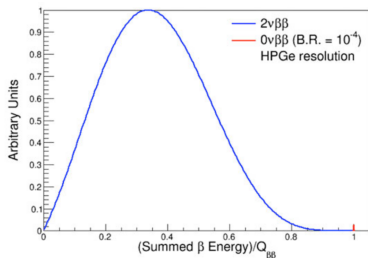


$\nu$  oscillation discovered by Kajita (Super-K) and McDonald (SNO)  
[Nobel prize 2015]

- New possibility for  $0\nu\beta\beta$  search  $\Rightarrow$  sensitive to neutrino's absolute mass

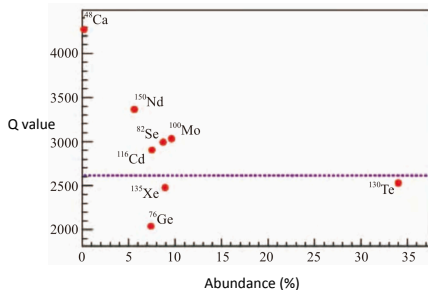
# Experimental search

## $0\nu\beta\beta$ vs $2\nu\beta\beta$ decay



$T_{1/2}^{0\nu} > 10^{26}$  yr  $\Rightarrow$  **Ton of isotopes  $\sim 10^{28}$  nuclei**

$\Rightarrow$  requires both large decay energy (Q value) and isotope abundance



# Experiments underway

## $0\nu\beta\beta$ decay

- The easiest way to determine whether  $\nu$  is a Majorana fermion
- Give the information on the absolute mass scale of  $\nu$
- Provide the evidence of lepton number violation

## More than 10 experiments underway

CANDLES	Ca-48	60 CaF <sub>2</sub> crystals in liq. scint	6 kg	Construction
CARVEL	Ca-48	<sup>48</sup> CaWO <sub>4</sub> crystal scint.	100 kg	
COBRA	Cd-116, Te-130	CdZnTe detectors	10 kg	R&D
CUROICINO	Te-130	TeO <sub>2</sub> Bolometer	11 kg	Operating
CUORE	Te-130	TeO <sub>2</sub> Bolometer	206 kg	Construction
DCBA	Nd-150	Nd foils & tracking chambers	20 kg	R&D
EXO200	Xe-136	Xe TPC	200 kg	Construction
EXO	Xe-136	Xe TPC	1-10t	R&D
GEM	Ge-76	Ge diodes in LN	1 t	
GERDA	Ge-76	Seg. and UnSeg. Ge in	35-40 kg	Construction

- 4 Exp. (Majorana, EXO, CUORE, GERDA) reached  $T_{1/2}^{0\nu} > 10^{25}$  year
- 1 Exp. (KamLAND-Zen) exceeded the level of  $1 \times 10^{26}$  year

## PandaX reports the lower limit of $T_{1/2}^{0\nu} > 2.1 \times 10^{23}$ from Chinese experiments

### PandaX 实验组发表首个寻找马约拉纳中微子结果

上海交通大学物理与天文学院 1周前



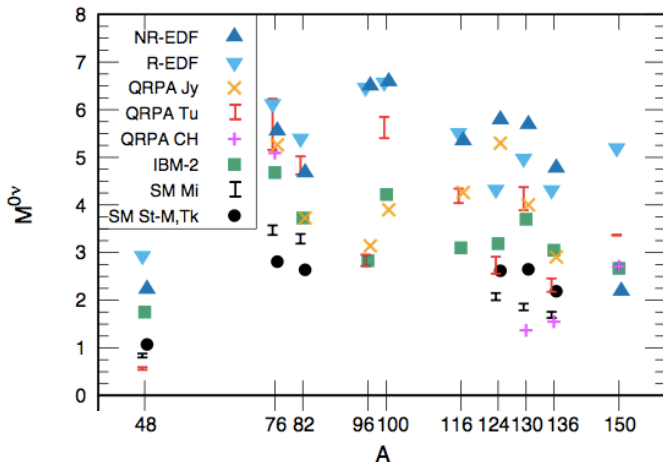
上海交通大学牵头的PandaX合作组近期在“中国物理C”杂志（IF=5.86）以“编辑推荐”的形式发表了首个利用液氙探测器寻找马约拉纳中微子的实验结果。他们的实验表明，核同位素 $^{136}\text{Xe}$ 通过马约拉纳中微子而产生的衰变寿命大于 $2.1 \times 10^{23}$ 年，即两千一百万亿亿年，比宇宙的年龄长了约15万亿倍。这是在中国本土物理实验中探测到最长的核素衰变寿命下限。该结果也能解释为马约拉纳电子中微子有效质量的上限， $m_{ee} < 1.4 \sim 3.7 \text{eV}/c^2$  (电子伏特)。

# Double $\beta$ decay: generic difficulties

At present, lattice QCD mainly targets on light nuclei

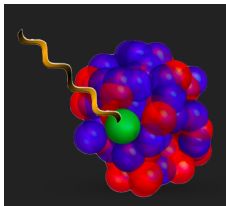
- For nucleus A:  $\frac{\text{signal}}{\text{noise}} \sim \exp[-A(M_N - 3/2m_\pi)t] \Rightarrow$  a sign problem!

For nuclear matrix element, various models yield O(100%) discrepancies

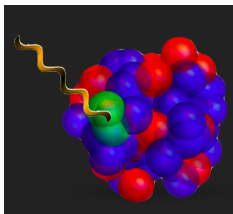


## Coupling of currents to nuclei in nuclear EFT [Detmold, talk at Lat18]

- One body coupling dominates



- Two nucleon contributions are subleading but non-negligible



A promising way to provide few-body inputs to ab initio many-body calculations

## Progress and Challenges in Neutrinoless Double Beta Decay

ECT\* workshop subscription



ECT\*, Strada delle Tabarelle, 286, Villazzano, 38123 Trento, Italy

Monday, 15 July 2019 at 08:00 - Friday, 19 July 2019 at 18:00 (CEST)



### Summarize on recent advances in

- Lattice QCD
- Chiral effective field theory
- Many-body nuclear theory

### Target on

- a seamless connection between the theory at quark and nuclear level
- reliable calculations of the nuclear matrix elements, with robust uncertainty



If neutrinoless double beta decays exist ...

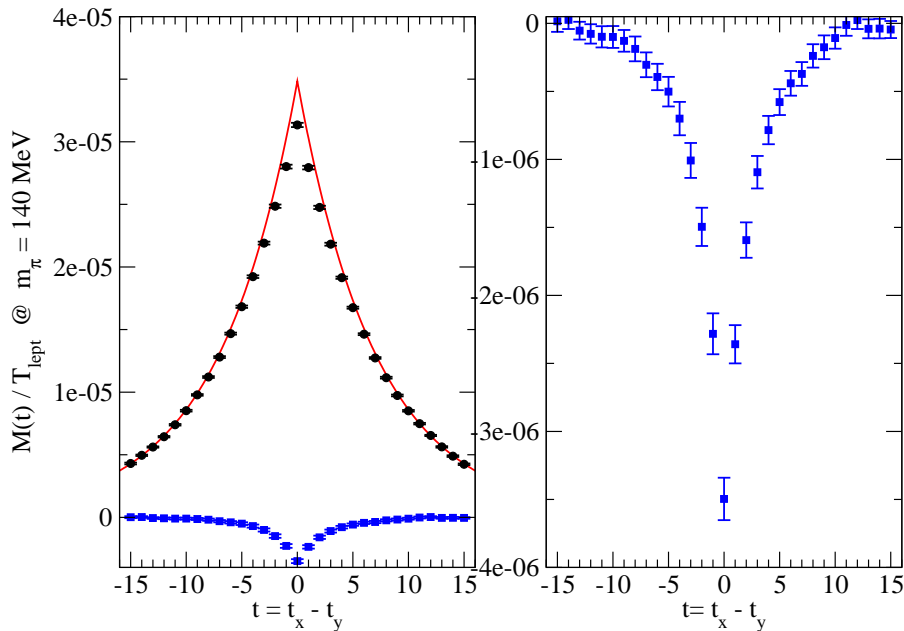
$$\pi^- \pi^- \rightarrow ee \text{ and } \pi^- \rightarrow \pi^+ ee$$

Lattice QCD starts with simplest decays

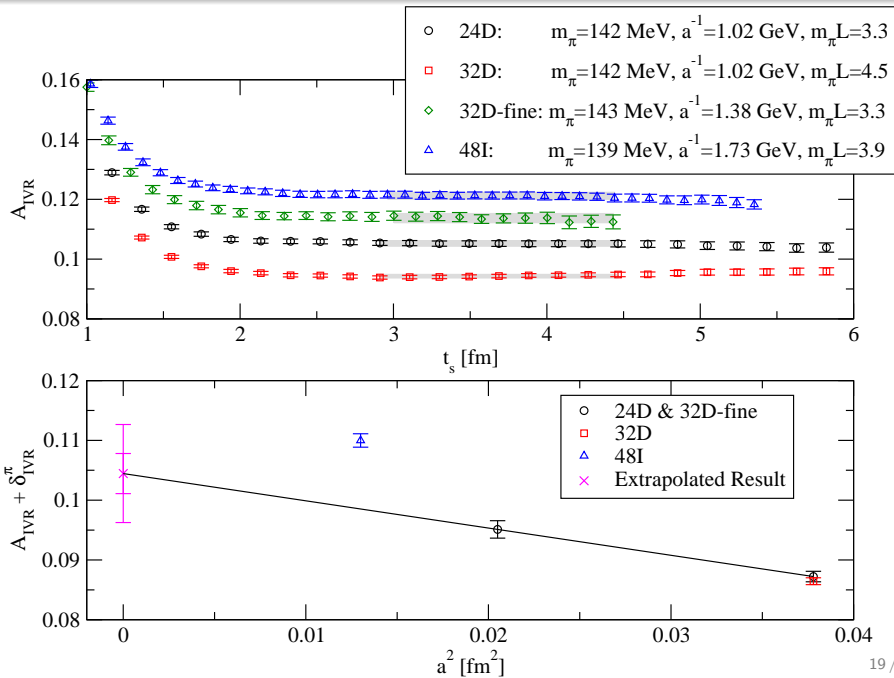
$\pi^- \pi^- \rightarrow ee$ : XF, L. Jin, X. Tuo, S. Xia, PRL122 (2019) 022001

$\pi^- \rightarrow \pi^+ ee$ : X. Tuo, XF, L. Jin, arXiv:1909.13525, accepted by PRD

# $\pi\pi \rightarrow ee$ decay amplitude @ $m_\pi = 140$ MeV



# $\pi^- \rightarrow \pi^+ ee$ : infinite volume reconstruction



# Summary of $\pi^- \pi^- \rightarrow ee$ and $\pi^- \rightarrow \pi^+ ee$

## Chiral perturbation theory for $\pi^- \pi^- \rightarrow ee$

[Cirigliano, Dekens, Mereghetti, Walker-Loud, PRC97 (2018) 065501]

$$\frac{\mathcal{A}(\pi^- \pi^- \rightarrow ee)}{2F_\pi^2 T_{\text{lept}}} = 1 - \frac{m_\pi^2}{(4\pi F_\pi)^2} \left( 3 \log \frac{\mu^2}{m_\pi^2} + \frac{7}{2} + \frac{\pi^2}{4} + \frac{5}{6} g_\nu^{\pi\pi}(\mu) \right)$$

## Lattice calculation yields (statistical error only)

[XF, L. Jin, X. Tuo, S. Xia, PRL122 (2019) 022001]

$$\frac{\mathcal{A}(\pi\pi \rightarrow ee)}{2F_\pi^2 T_{\text{lept}}} = 0.910(3) \quad \Rightarrow \quad g_\nu^{\pi\pi}(m_\rho) = -12.0(3)$$

## Chiral perturbation theory for $\pi^- \rightarrow \pi^+ ee$

$$\frac{\mathcal{A}(\pi^- \rightarrow \pi^+ ee)}{2F_\pi^2 T_{\text{lept}}} = 1 + \frac{m_\pi^2}{(4\pi F_\pi)^2} \left( 3 \log \frac{\mu^2}{m_\pi^2} + 6 + \frac{5}{6} g_\nu^{\pi\pi}(\mu) \right)$$

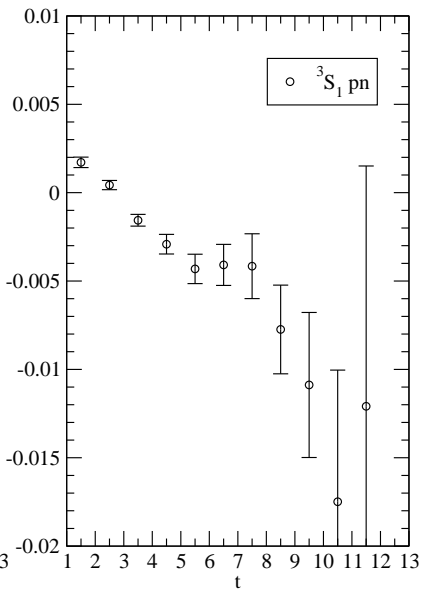
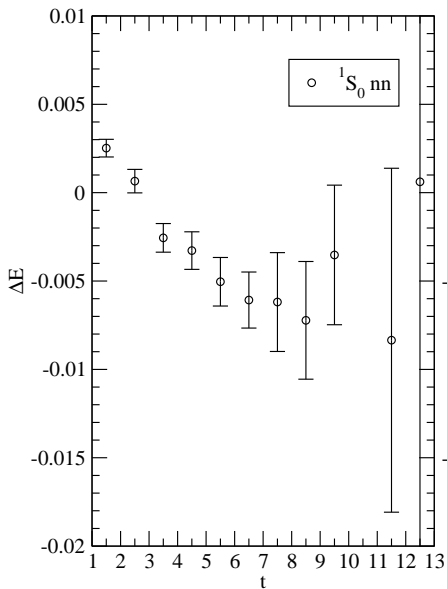
## Lattice calculation yields (statistical + systematical errors)

[X. Tuo, XF, L. Jin, accepted by PRD]

$$\frac{\mathcal{A}(\pi^- \rightarrow \pi^+ ee)}{2F_\pi^2 T_{\text{lept}}} = 1.105(3)(7) \quad \Rightarrow \quad g_\nu^{\pi\pi}(m_\rho) = -10.9(3)(7)$$

Move to dibaryon

# Di-neutron vs deuteron



## 2 point correlation function

- Bounding energy

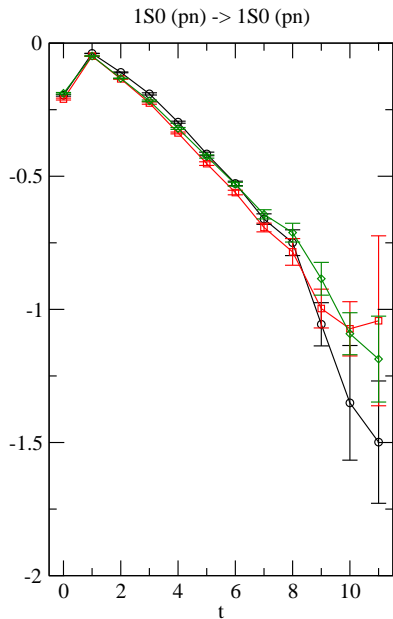
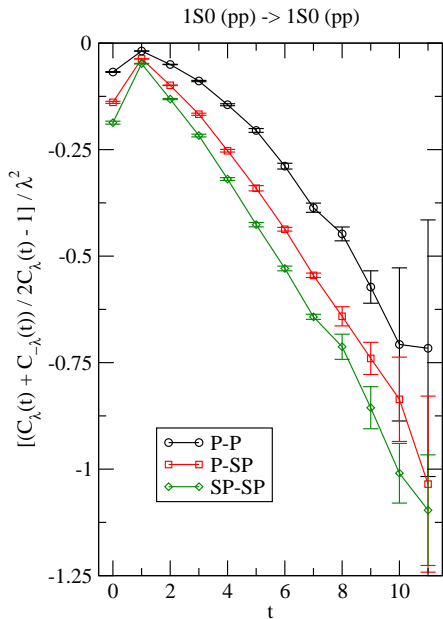
## 3 point correlation function

- $g_A$  quenching effects
- proton-proton fusion  $pp \rightarrow de^+\nu_e$

## 4 point correlation function

- $2\nu 2\beta$  decay:  $nn \rightarrow ppee\bar{\nu}\bar{\nu}$
- $0\nu 2\beta$  decay:  $nn \rightarrow ppee$

# 4 point correlation function for $0\nu 2\beta$ decay





## With developments of supercomputers, computational technologies and novel ideas and methods

- the precision of lattice QCD calculation has been improved significantly and have great potential to further improve

## For flavor physics:

- lattice QCD provides useful low-energy QCD information
- plays important role in high-precision frontier

## The techniques developed in flavor physics can be used in nuclear physics

- help to study the rare processes related to nuclear matter
- Can one day, nuclear physics become a new flavor physics?



# Double $\beta$ decay of nuclei

Cirigliano, Dekens, Mereghetti, Walker-loud, PRC97 (2018) 065501

Begin with the effective Lagrangian  $\mathcal{L}_{\text{eff}}$  for the single  $\beta$  decay

$$\mathcal{L}_{\text{eff}} = 2\sqrt{2}G_F V_{ud}(\bar{u}_L\gamma_\mu d_L)(\bar{e}_L\gamma_\mu\nu_{eL})$$

Contributions are identified into three regions in EFT

- Hard region:  $\Lambda \gg 1$  GeV

$$\int d^4x e^{i\Lambda x} \mathcal{L}_{\text{eff}}(x)\mathcal{L}_{\text{eff}}(0) \sim 8G_F^2 V_{ud}^2 \frac{m_{\beta\beta}}{\Lambda^2} (\bar{u}_L\gamma_\mu d_L)(\bar{u}_L\gamma_\mu d_L)\bar{e}_L e_L^c.$$

In lattice QCD, a hard cutoff is introduced by  $1/a \Rightarrow O(a^2)$  effects

- Soft region:  $O(100$  MeV) -  $O(1$  GeV)
  - ▶ Few-body decay dominates
  - ▶ Nuclear potential mediated by pions:  $\pi\pi \rightarrow ee$ ,  $\pi n \rightarrow pee$ ,  
 $nn \rightarrow ppee$ ,  $\dots$
- Ultrasoft or radiative region:  $\Lambda \ll 100$  MeV
  - ▶ Neutrinos feel the complete nucleus instead of just the nucleons

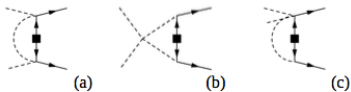
# Diagrams in chiral perturbation theory

## LO diagrams (tree level)

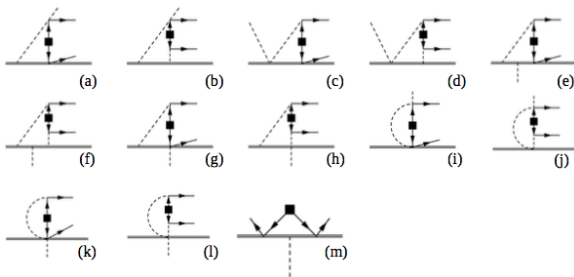


## NLO diagrams (one loop)

•  $\pi\pi \rightarrow ee$



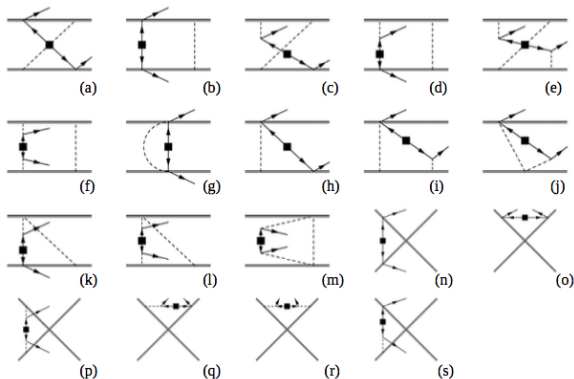
•  $\pi n \rightarrow pee$



# Diagrams in chiral perturbation theory

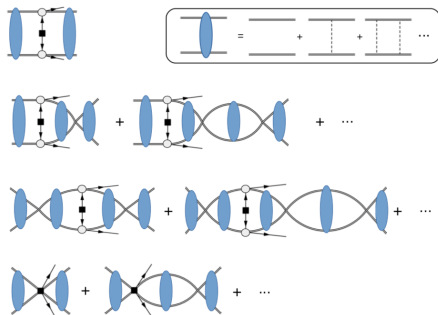
## NLO diagrams (one loop)

•  $nn \rightarrow ppee$



# Logarithmic divergence at NNLO

Cirigliano, Dekens, De Vries, et.al. PRL120 (2018) 202001



## Transition amplitude from EFT

$$A_\nu = \langle pp, \text{out} | V_\nu | nn, \text{in} \rangle = - \int d^3\mathbf{r} \psi_{\mathbf{p}'}^-(\mathbf{r})^* V_\nu(\mathbf{r}) \psi_{\mathbf{p}}^+(\mathbf{r})$$

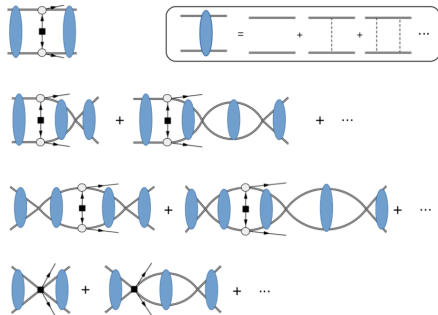
## Wave function

$$\psi_{\mathbf{p}}^\pm(\mathbf{r}) = \chi_{\mathbf{p}}^\pm(\mathbf{r}) + \chi_{\mathbf{p}}^\pm(\mathbf{0}) K_E G_E^\pm(\mathbf{r}, \mathbf{0})$$

where

$$G_E^\pm(\mathbf{r}, \mathbf{r}') = \langle \mathbf{r} | \frac{1}{E - H_0 - V_\pi \pm i\epsilon} | \mathbf{r}' \rangle, \quad K_E = \frac{C}{1 - CG_E^+ \mathbf{0}, \mathbf{0}}$$

# Logarithmic divergence at NNLO



## Transition amplitude from EFT

$$A_\nu = \langle pp, \text{out} | V_\nu | nn, \text{in} \rangle = - \int d^3\mathbf{r} \psi_{\mathbf{p}'}^-(\mathbf{r})^* V_\nu(\mathbf{r}) \psi_{\mathbf{p}}^+(\mathbf{r})$$

## For $r \rightarrow 0$

$V_\nu(\mathbf{r}) \rightarrow 1/r$ , Neutrino potential

$\chi_{\mathbf{p}}^+(\mathbf{r}) \rightarrow \text{constant}$ , Yukawa wave function

$G_E^\pm(\mathbf{r}, \mathbf{0}) \rightarrow \frac{m_N}{4\pi r} + \dots$ , Propagator

Need lattice QCD to confirm it

# How much chance to observe $0\nu\beta\beta$ ?

- The two mass eigenstates that mix most strongly with electron flavor are lighter than the third (**normal hierarchy**) or heavier (**inverted hierarchy**)

