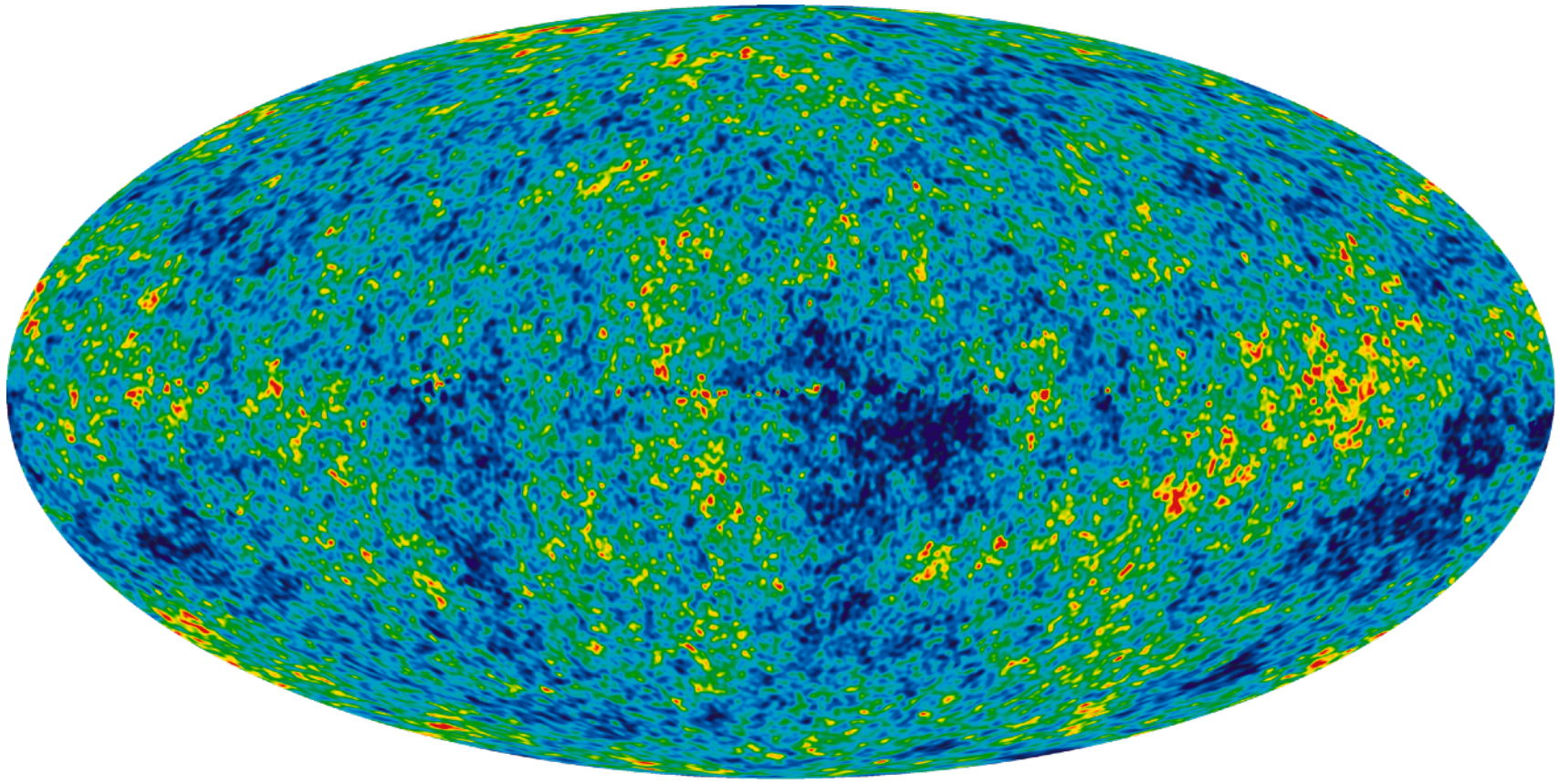


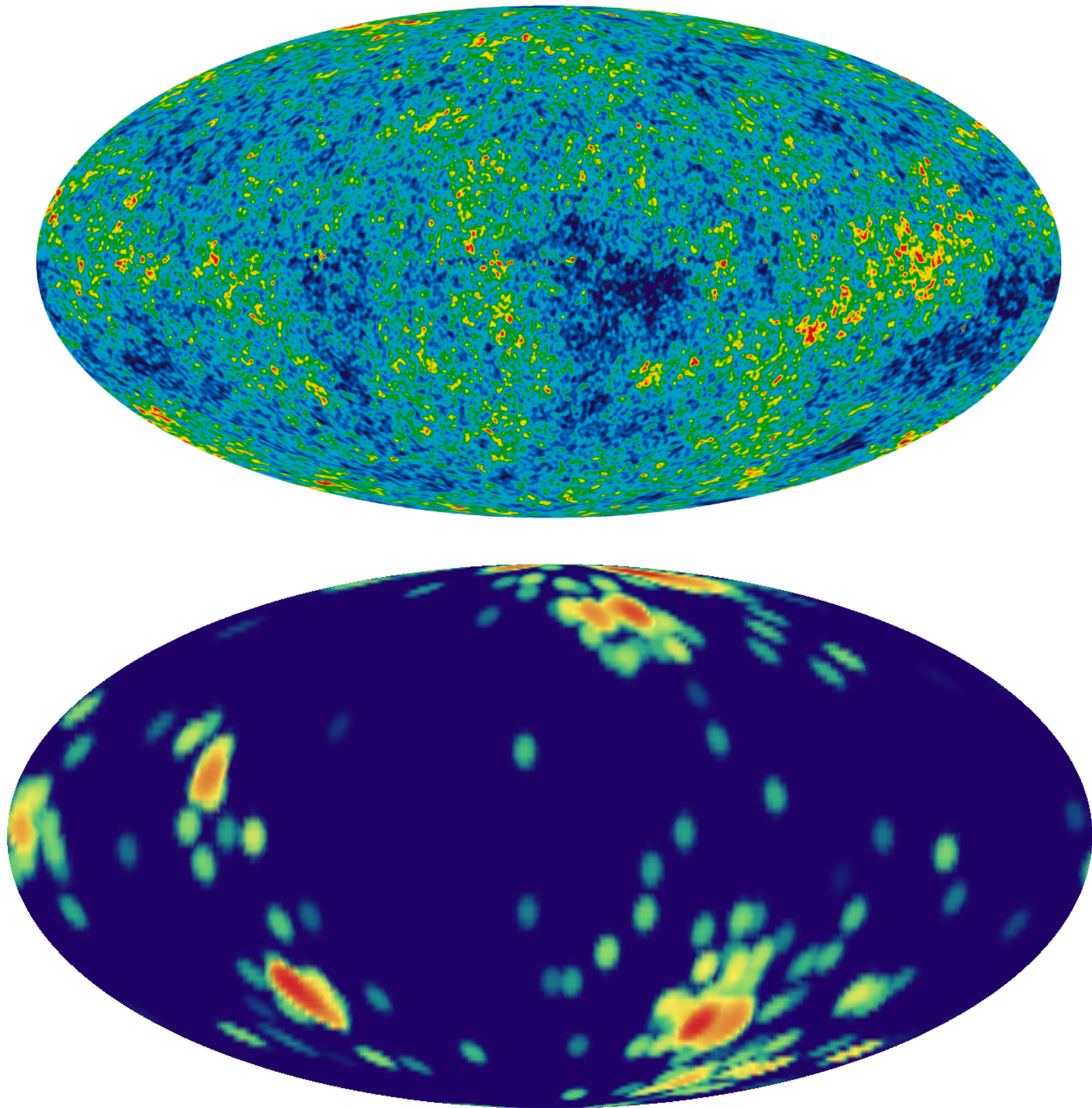
Three-Point Energy Correlator in the Coplanar Limit

Hua Xing Zhu
Zhejiang University

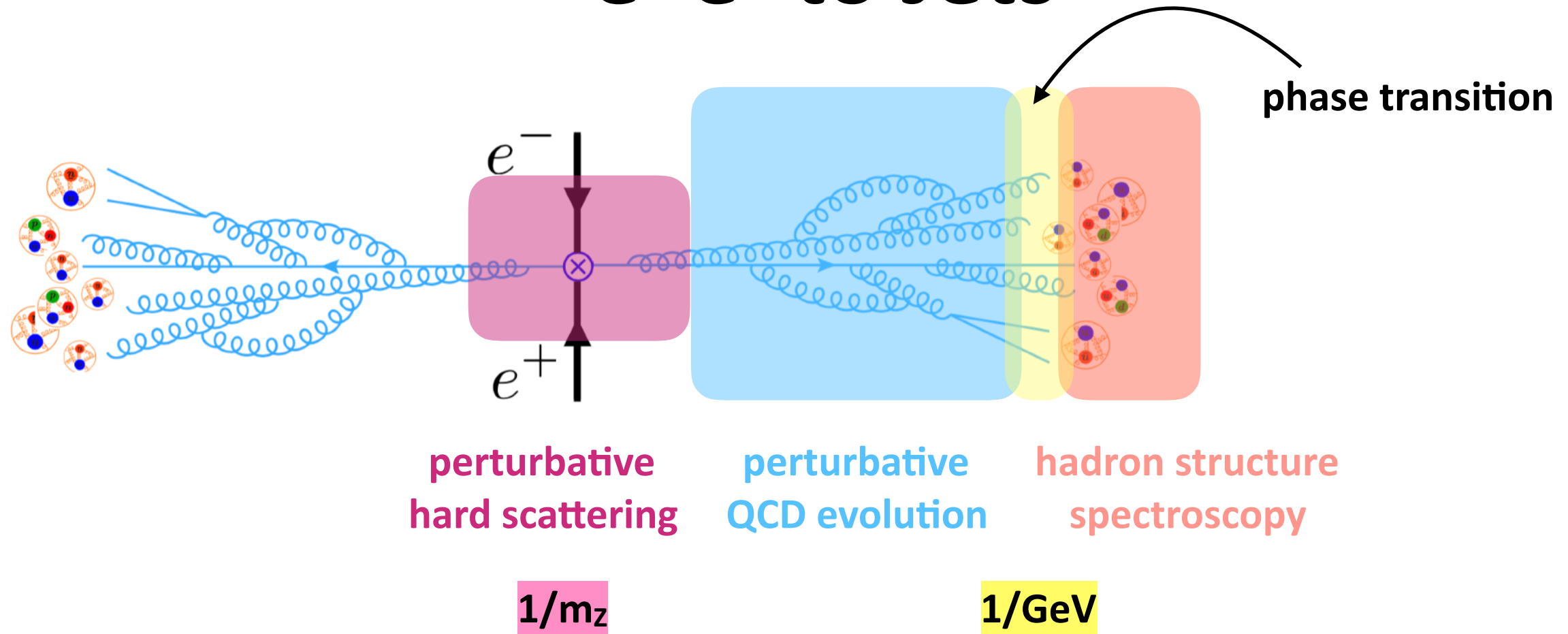
Work in preparation with Anjie Gao, Ming-Xing Luo, Tong-Zhi Yang

Beijing CEPC International Workshop
Nov. 18-20, 2019, IHEP

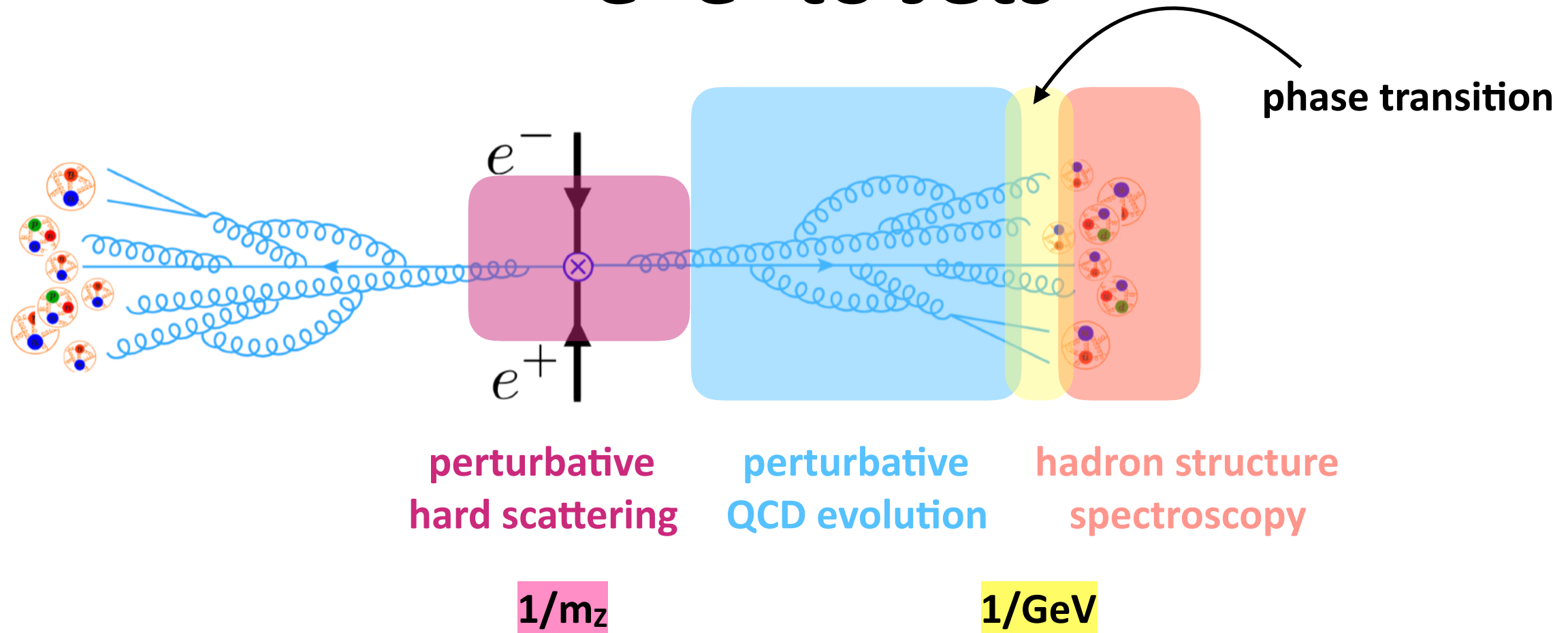




e^+e^- to Jets

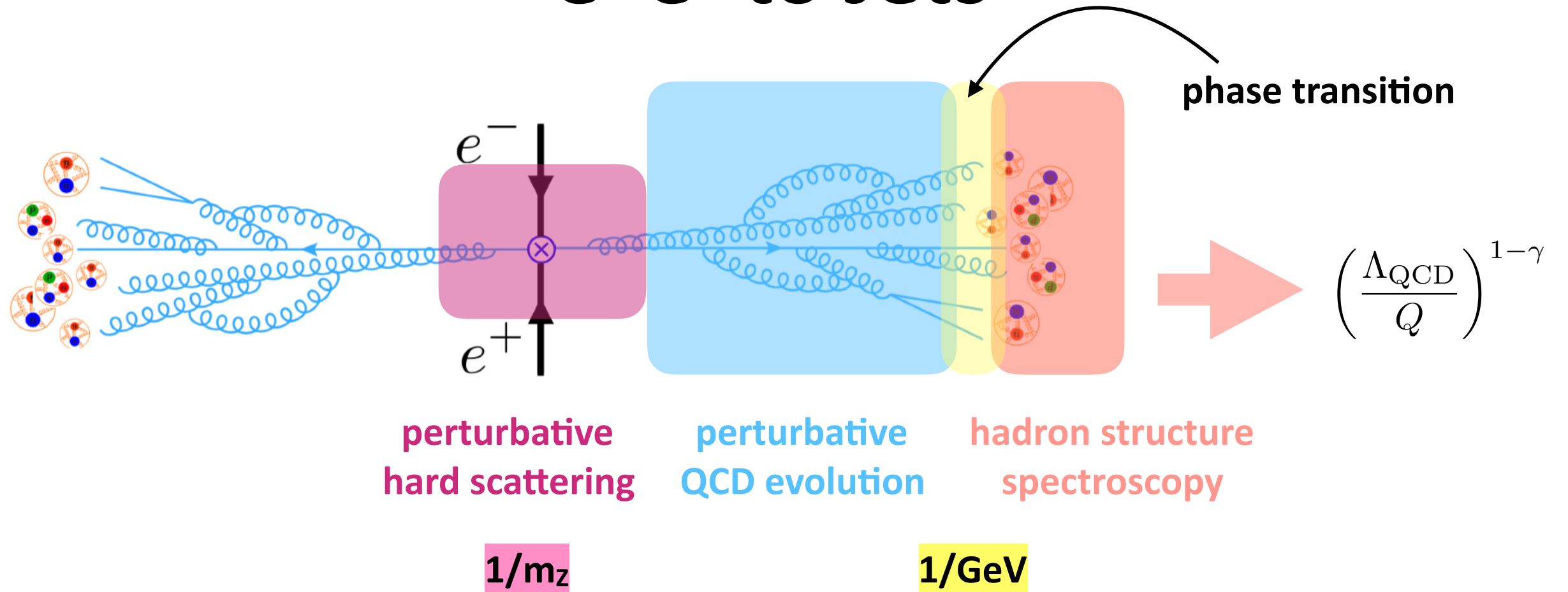


e^+e^- to Jets

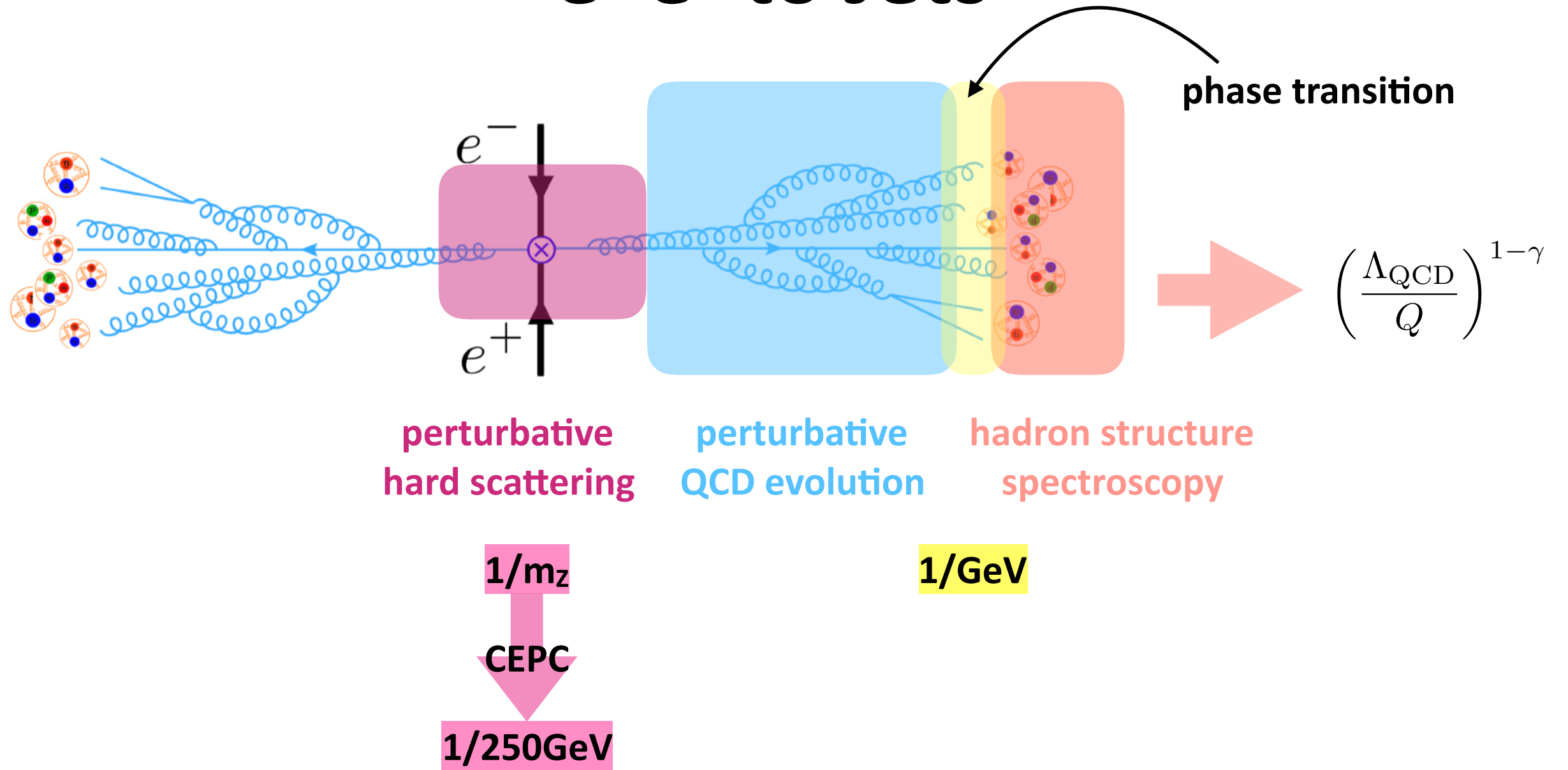


But isn't all of these physics can be studied at LEP?

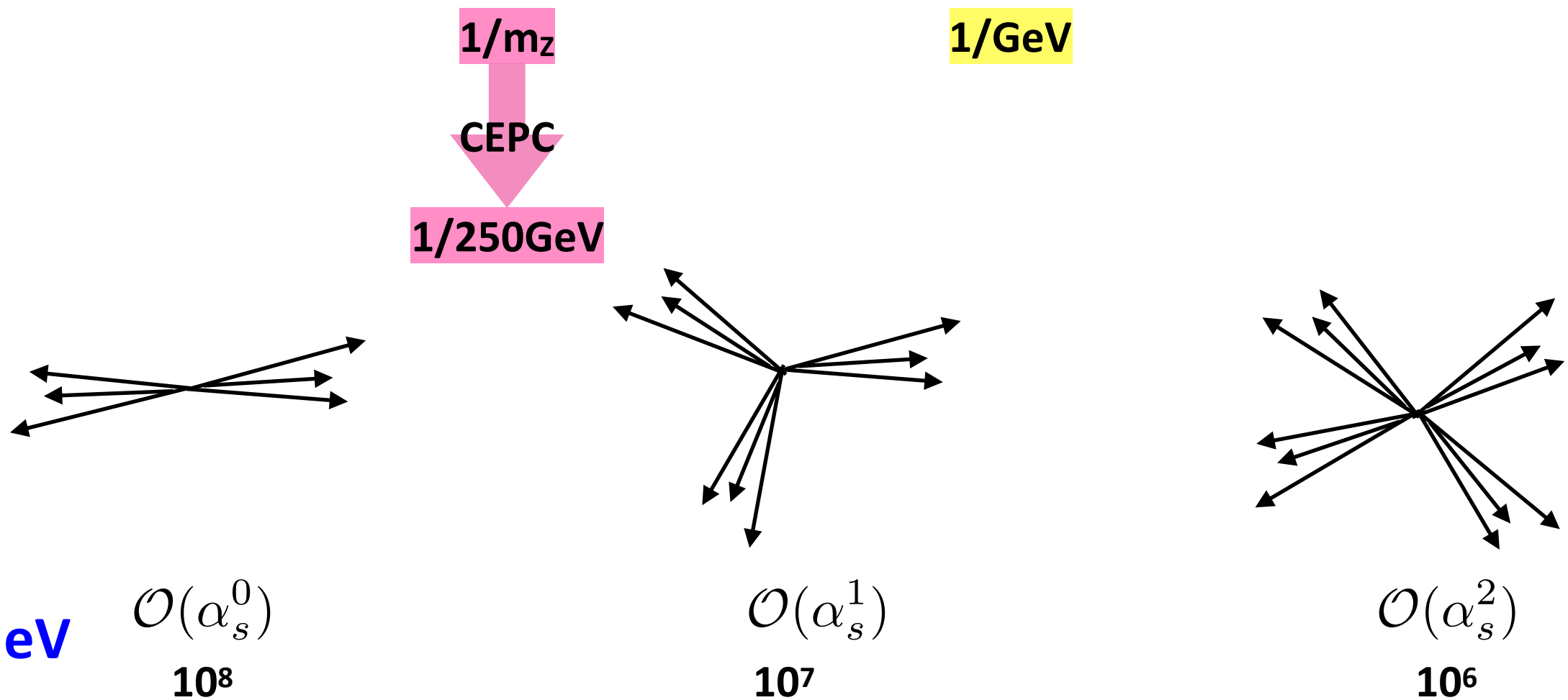
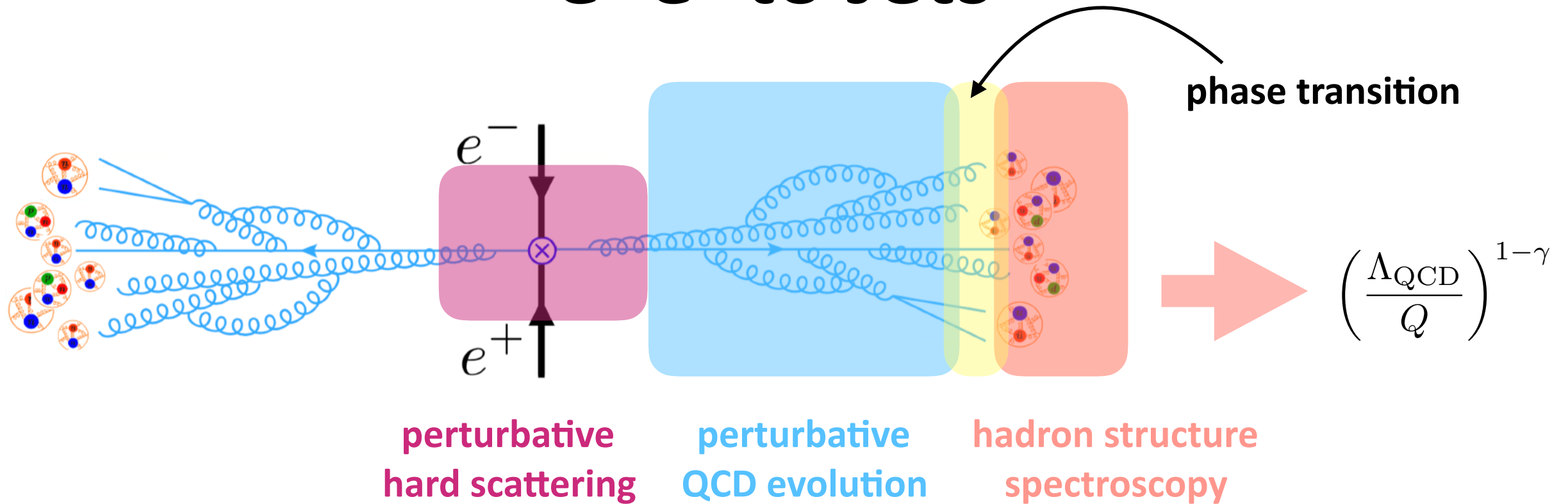
e+e- to Jets



e+e- to Jets



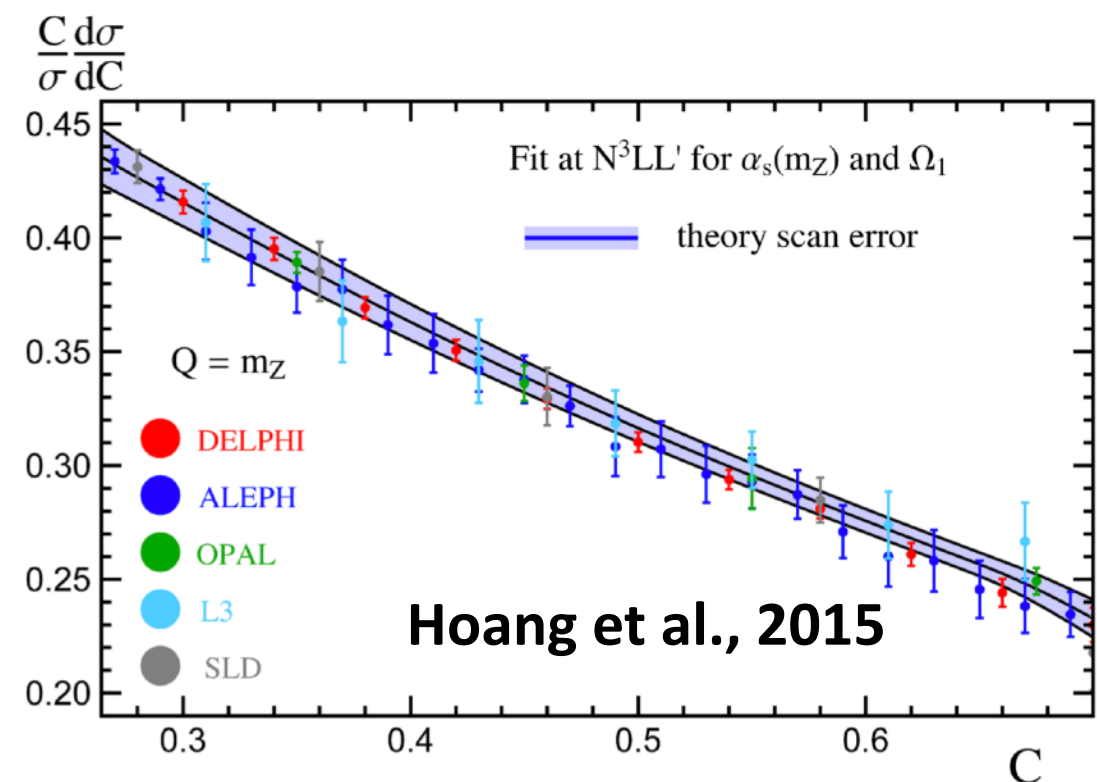
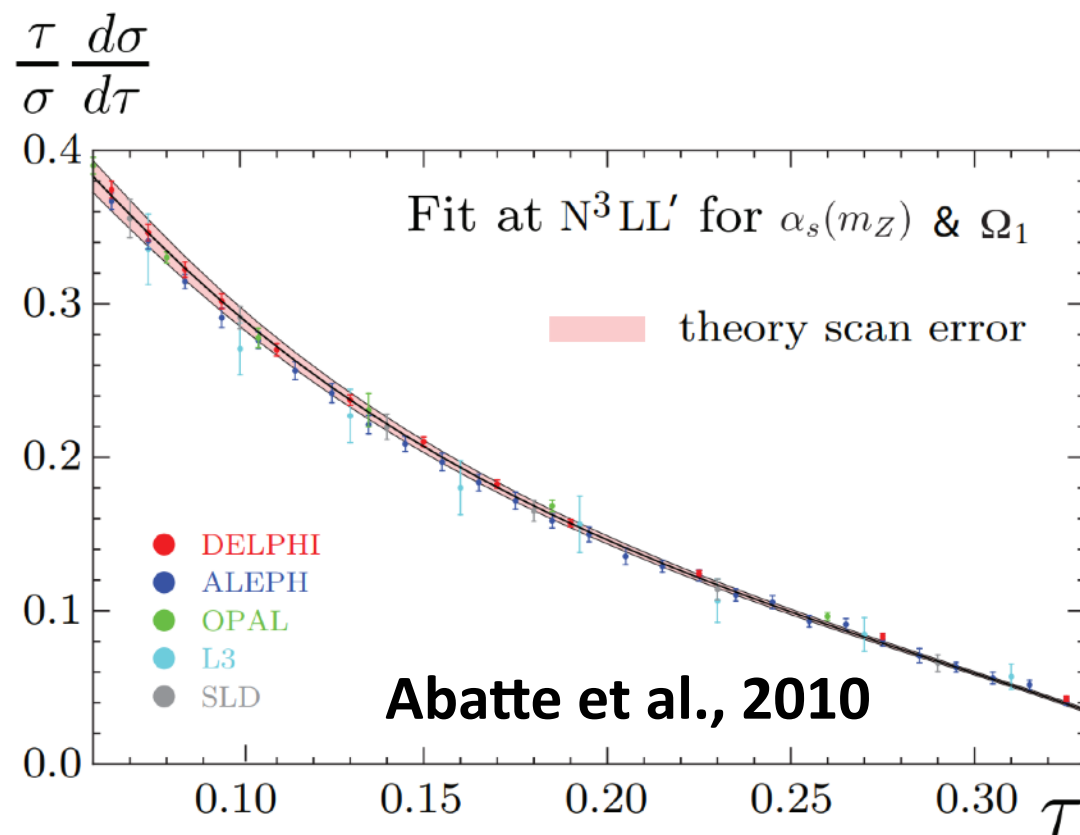
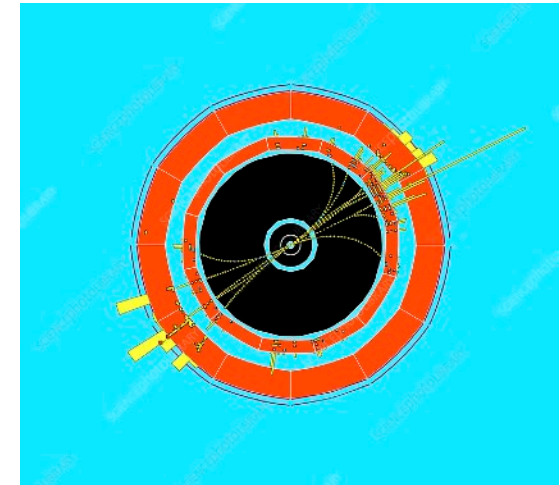
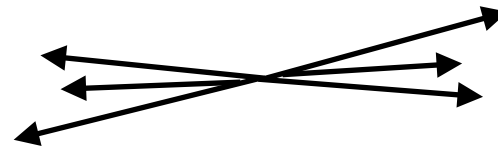
e+e- to Jets



Dijet Event Shape Observables

Energy flow observable which vanish in the dijet limit

Example: Thrust, C parameter, wide and total jet broadening, heavy jet mass, et al.

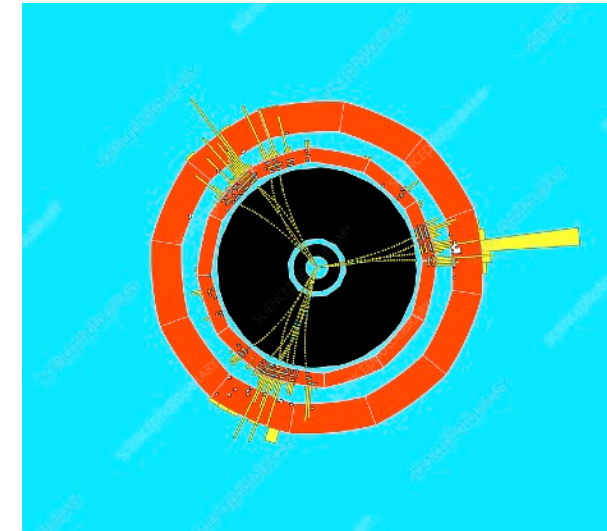
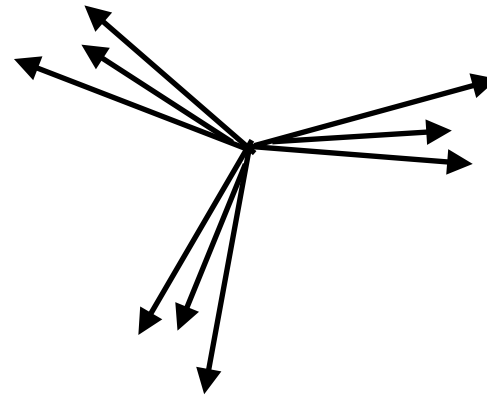


$$\alpha_s(m_Z) = 0.1135 \pm (0.0002)_{\text{exp}} \pm (0.0005)_{\text{hadr}} \pm (0.0009)_{\text{pert}}$$

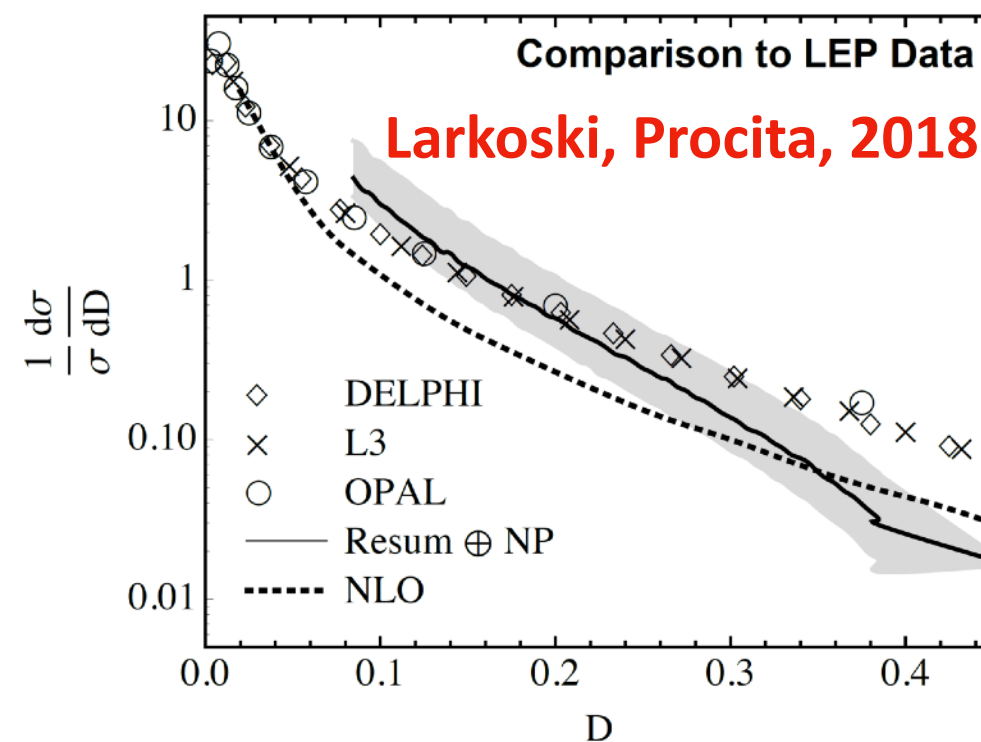
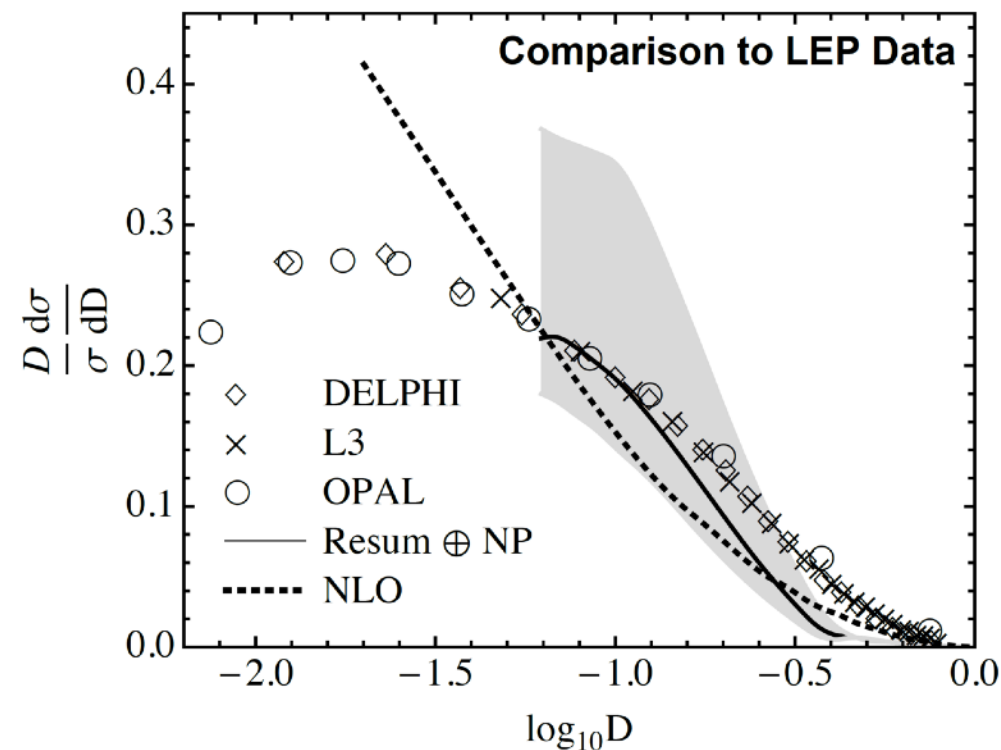
See also Andrii Verbytskyi's talk

Three Jet Event Shape Observables

Energy flow observable which vanish in the three jet limit



- Expect to be more sensitive to α_s
- Much less studied in the literature before
- Thrust minor [Banfi et al., 2001]; D parameter [Banfi et al., 2001; Larkoski, Procita, 2018]. **All limited to NLL study**



Three-Point Energy Correlator

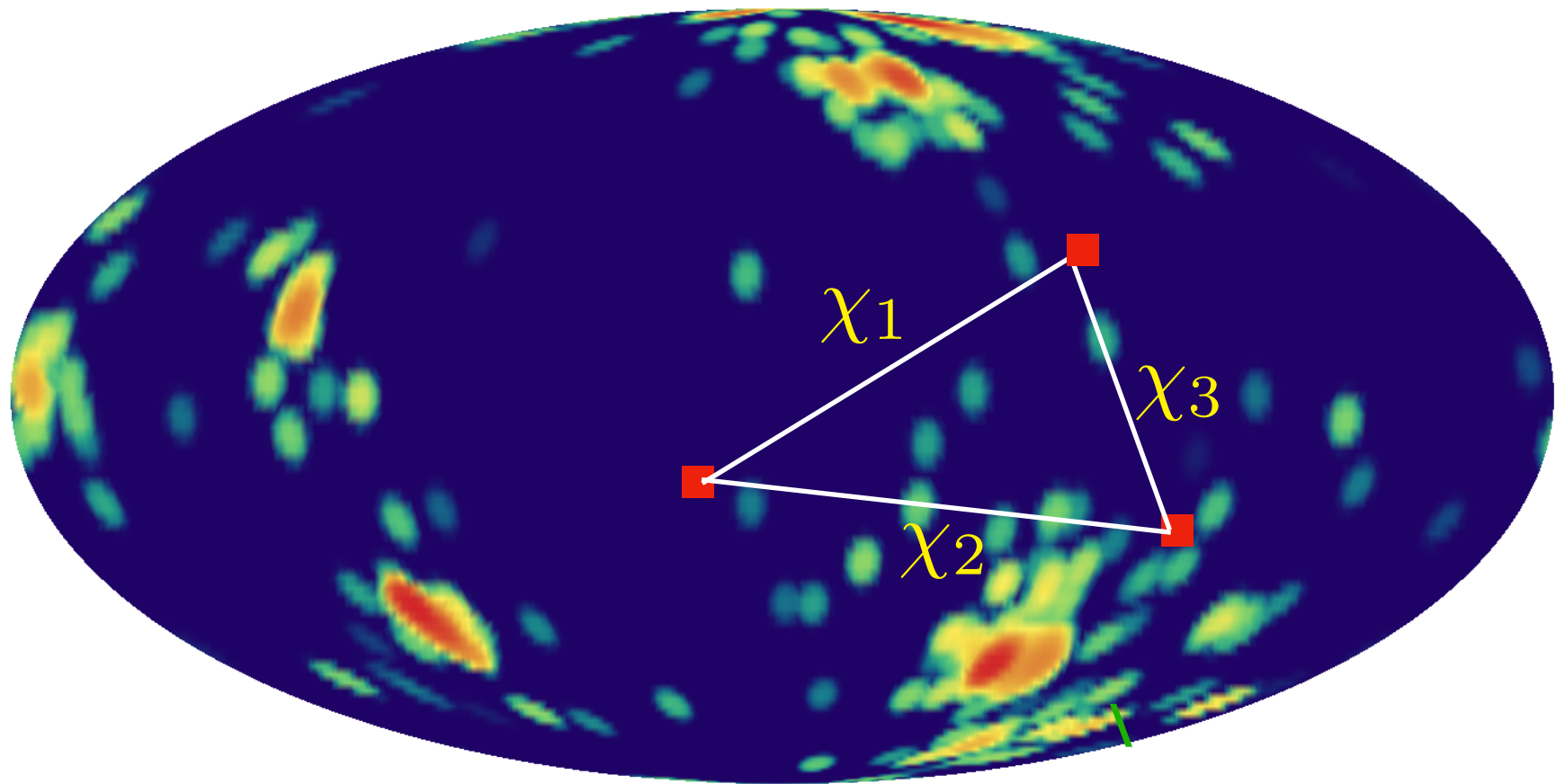
$$\sum_{ijk} \int d\sigma \frac{E_i E_j E_k}{Q^3 \sigma_{\text{tot}}} \delta(\cos \chi_1 - \cos \theta_{ij}) \delta(\cos \chi_2 - \cos \theta_{ik}) \delta(\cos \chi_3 - \cos \theta_{jk})$$

Σ

Sum
over all
events

Σ

Sum over all
inequivalent
configuration,
keeping $\chi_1, \chi_2,$
 χ_3 fixed



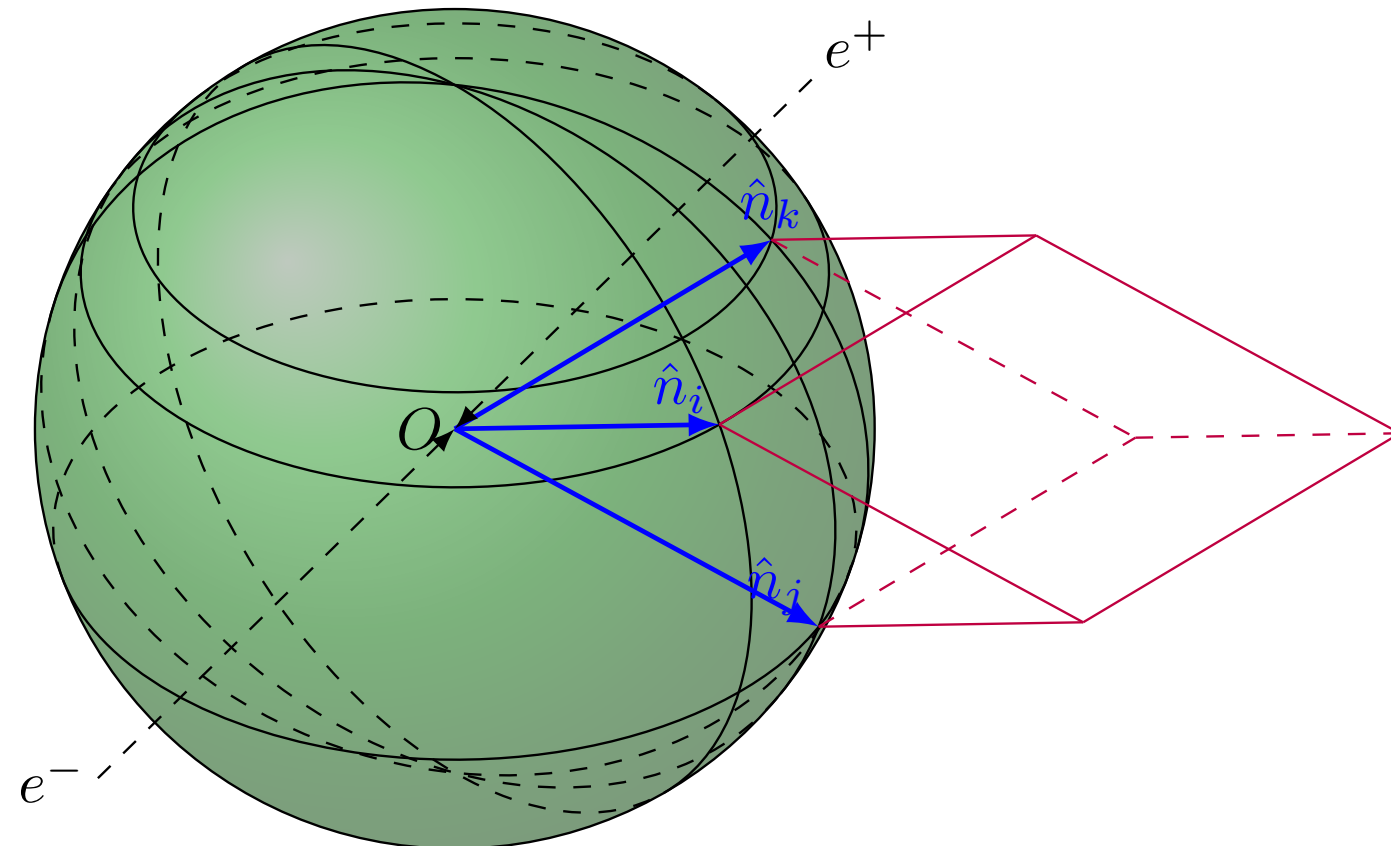
- Natural generalization of the Energy-Energy Correlator [Basham, Brown, Ellis, Love, 1978]
- Function of three angle χ_1, χ_2, χ_3

Energy Triple-Product Correlator

- The three angle χ_1, χ_2, χ_3 can also be specified by three unit vector $\vec{n}_1, \vec{n}_2, \vec{n}_3$
- The Energy Triple-Production Correlator (ETPC):

$$\tau_p = \left| \vec{n}_i \cdot (\vec{n}_j \times \vec{n}_k) \right|$$

$$\tau_p = \text{Vol}_{\text{hexahedron}}(\vec{n}_1, \vec{n}_2, \vec{n}_3)$$



Relation to D Parameter

- The D parameter is the product of the three eigenvalues of the sphericity tensor [Parisi, 1978; Donoghue et al., 1979]

$$\Theta_{\alpha\beta} = \frac{1}{Q} \sum_i \frac{p_{i\alpha} p_{i\beta}}{E_i}$$

$$D = 27\lambda_1\lambda_2\lambda_3$$

$$= \frac{27}{6} \{ (\lambda_1 + \lambda_2 + \lambda_3) [(\lambda_1 + \lambda_2 + \lambda_3)^2 - 3(\lambda_1^2 + \lambda_2^2 + \lambda_3^2)] + 2(\lambda_1^3 + \lambda_2^3 + \lambda_3^3) \}$$

$$= \frac{27}{6} \{ \text{Tr}\Theta [(\text{Tr}\Theta)^2 - 3\text{Tr}\Theta^2] + 2\text{Tr}\Theta^3 \}$$

$$= \frac{27}{Q^3} \sum_{i<j<k} \frac{|(\vec{p}_i \times \vec{p}_j) \cdot \vec{p}_k|^2}{E_i E_j E_k}$$

$$= \frac{27}{Q^3} \sum_{i<j<k} E_i E_j E_k \tau_{ijk}^2$$

Massless limit

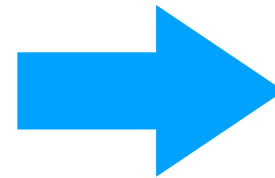
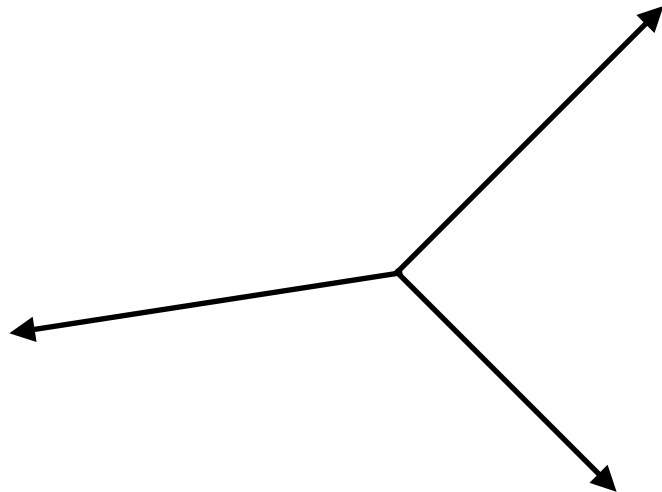
$$\langle D \rangle = \frac{9}{2} \int d\tau_p \tau_p^2 \text{ETPC}(\tau_p)$$

**Average of
D parameter**

Third moment of ETPC

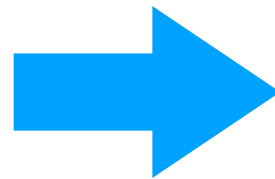
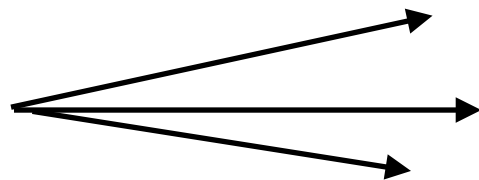
ETPC and Three Jet Configuration

3 jet coplanar
limit



$$\tau_p \rightarrow 0$$

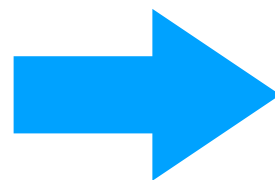
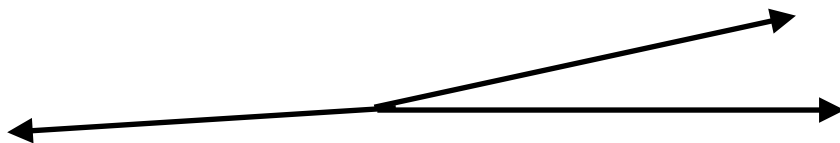
Triple collinear
limit



$$\tau_p \rightarrow 0$$

See also Ian
Moult's talk

Dijet double
collinear limit



$$\tau_p \rightarrow 0$$

We are interested in the 3 jet coplanar limit only in this talk

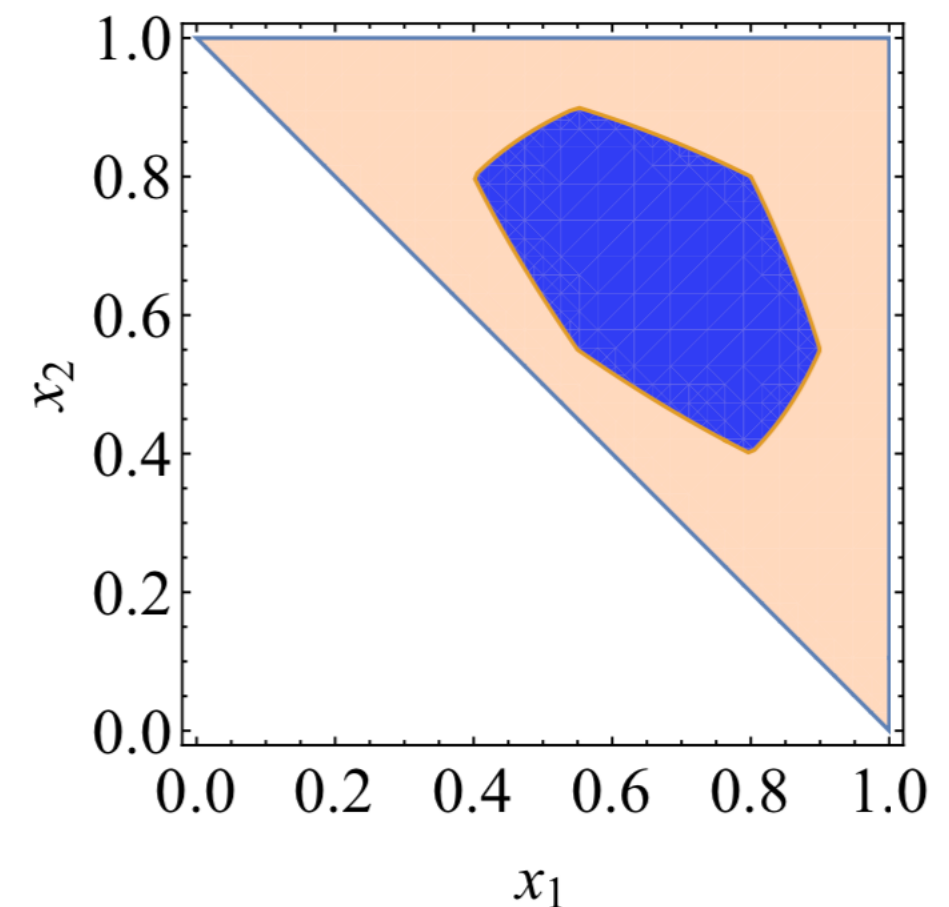
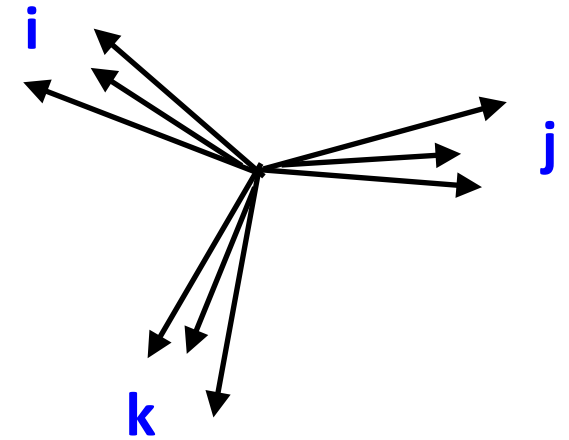
Phase Space Cut for 3 Jet Configuration

Use the k_T algorithm to find three jets, we keep the event only if $y_3 > y_{\text{cut}}$, and modify the definition of the ETPC to

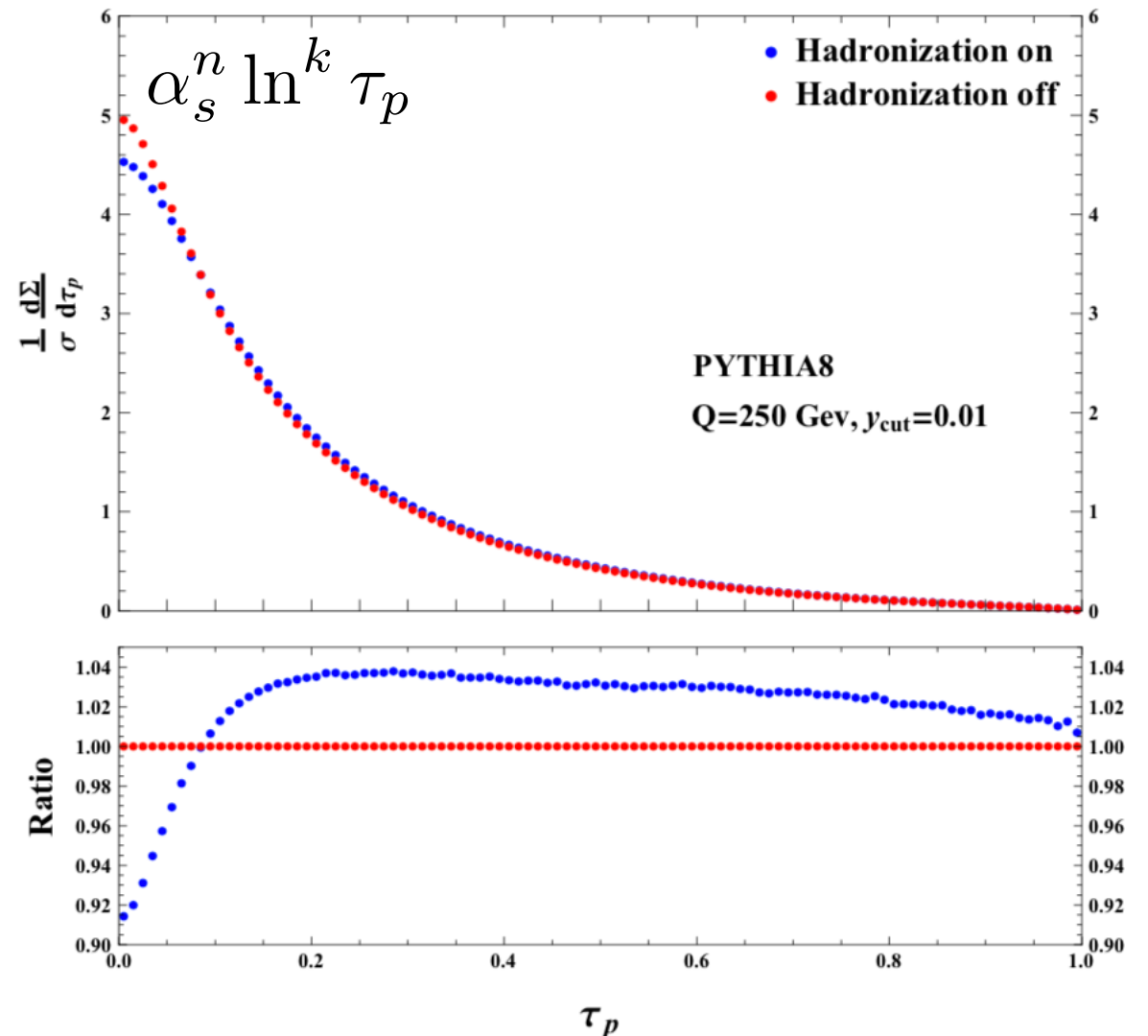
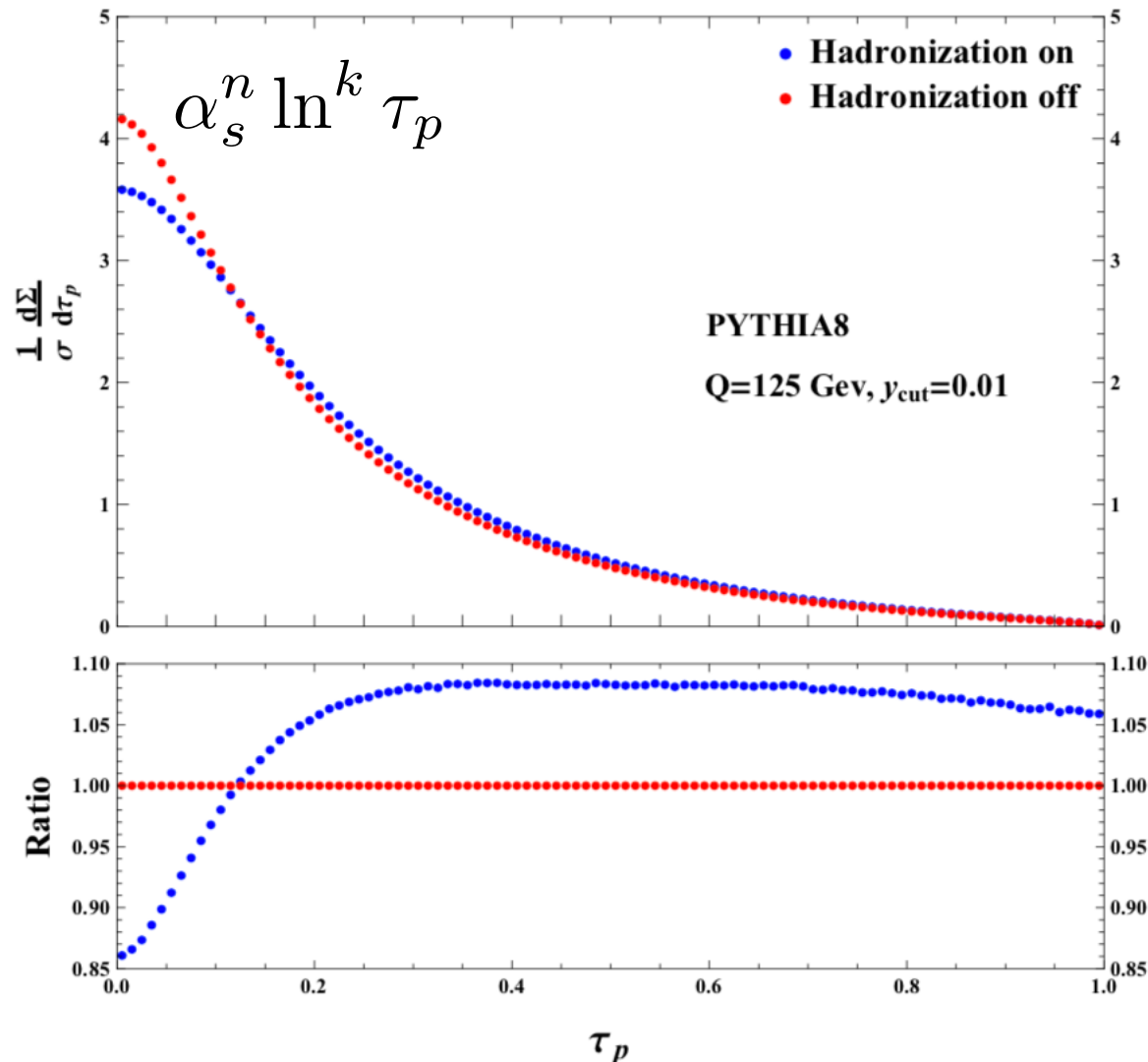
$$\sum_{\substack{i \in J_1 \\ j \in J_2 \\ k \in J_3}} \int d\sigma \frac{E_i E_j E_k}{E_{J_1} E_{J_2} E_{J_3} \sigma_{\text{tot}}} \delta(\tau_p - \tau_{ijk})$$

where J_1, J_2, J_3 denote three jets.

The phase space of three particles, with $y_{\text{cut}} = 0.1$
Here $x_i = 2E_i/Q$



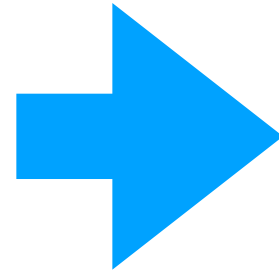
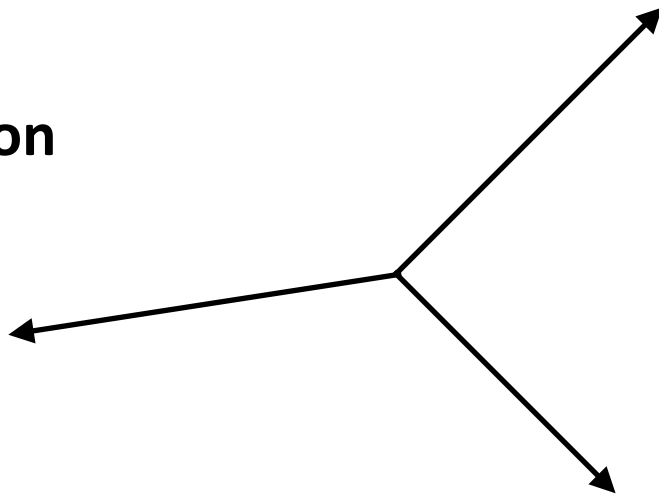
Pythia simulation for ETPC



- Substantial reduction of hadronization corrections from Z pole to 250 GeV
- Pythia simulation based on $ee \rightarrow qq$ configuration. No guarantee to get the normalization and shape right
- A first principle perturbative QCD calculation is needed

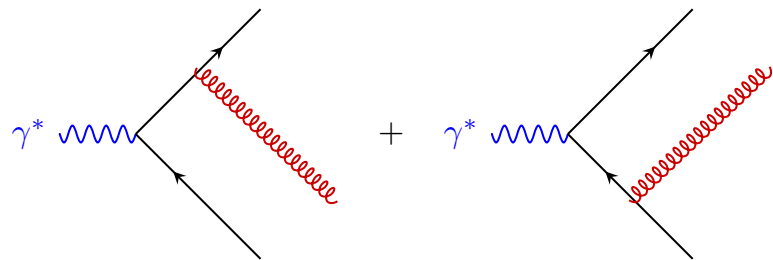
Short Distance Contribution

Born level
3 jet production

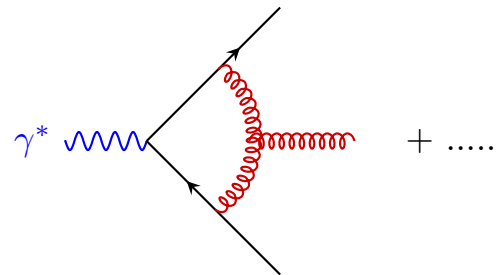


$$\tau_p \equiv 0$$

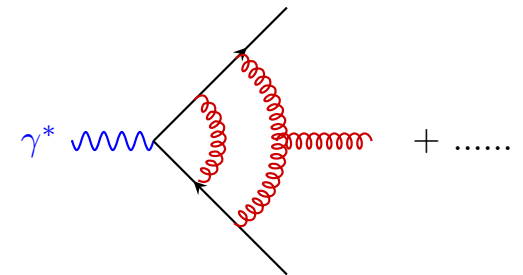
- Short distance corrections from $O(Q)$ hard virtual gluon do not change τ_p



LO



NLO

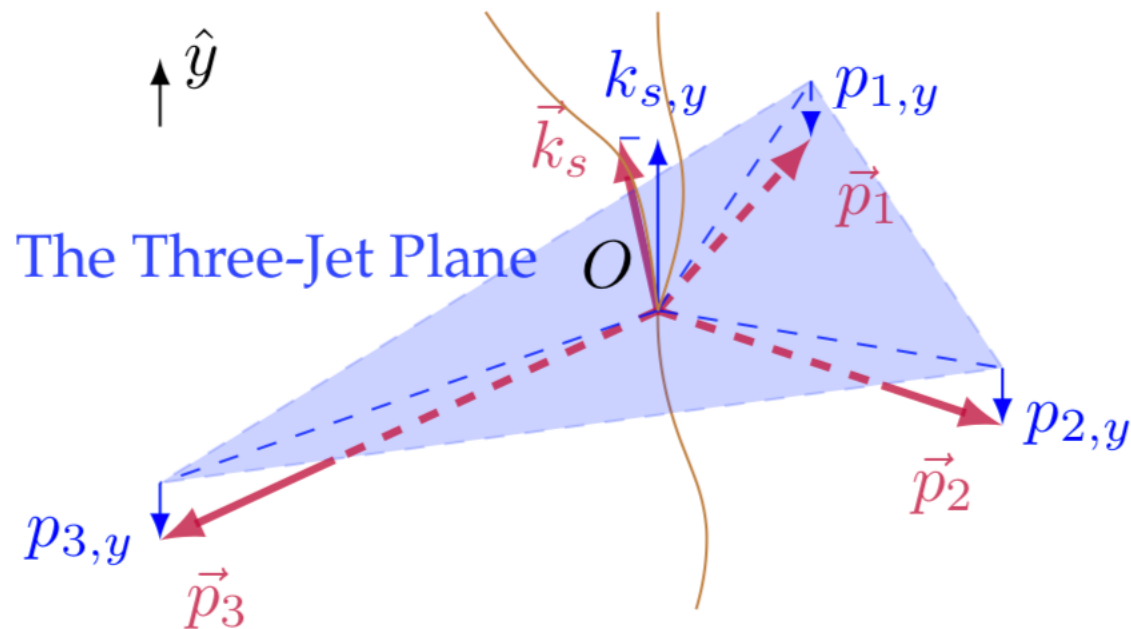


NNLO

- Higher order corrections are important to get the normalization right

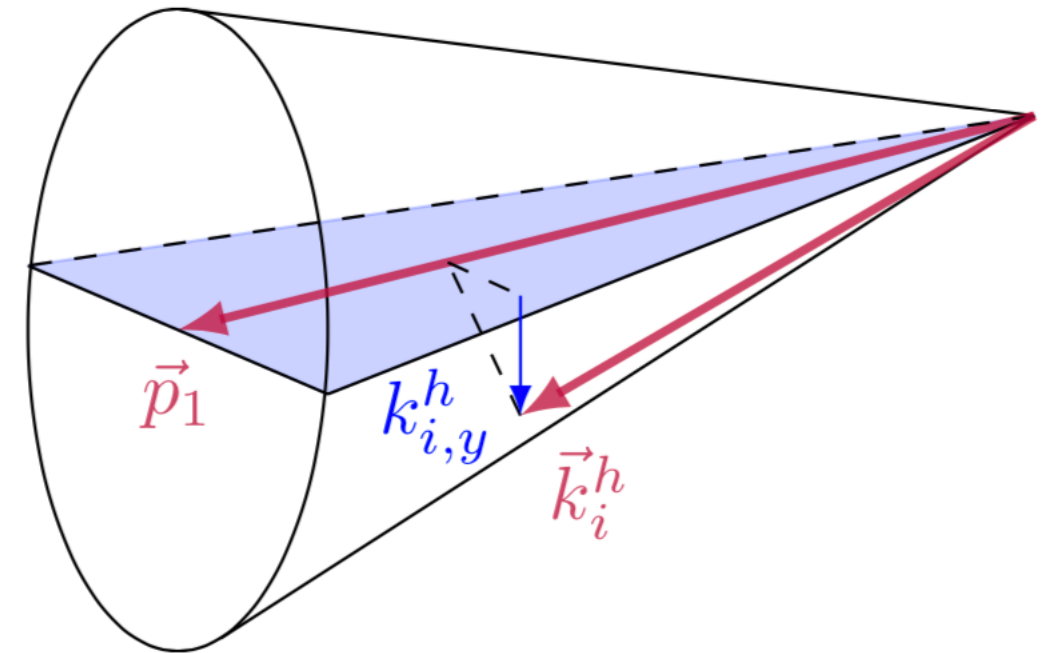
Soft/Collinear Contributions

- The recoil effect of soft radiations.



- The soft radiations *don't change the energy* of three jets, only make the three jets deviate the trijet plane slightly in the *opposite* direction.

- Collinear fragmentation

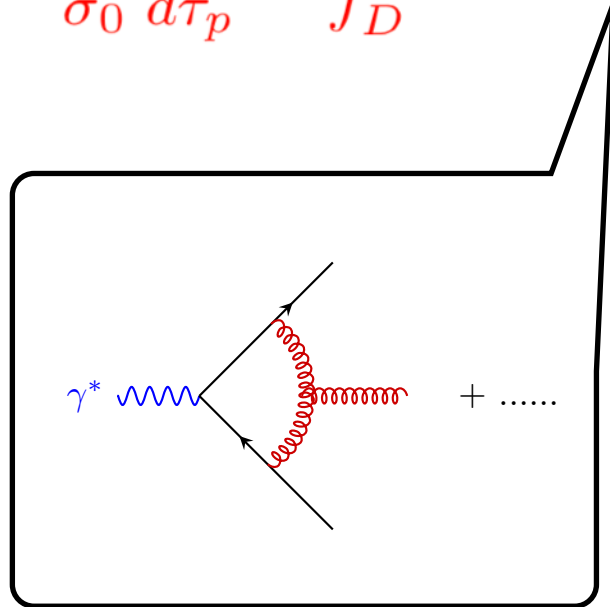


- The final state hardron carries longitudinal fraction momentum z_i^h from it's parent jet, i.e. $\vec{k}_i^h = z_i^h \vec{p}_i$.

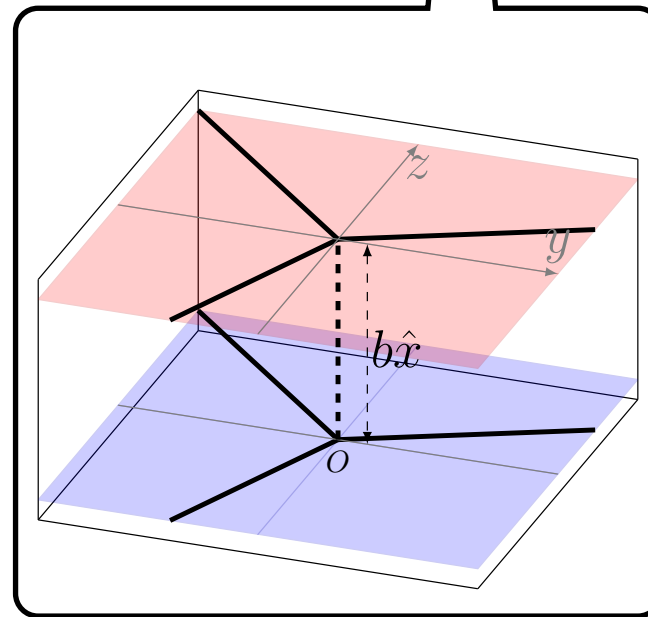
$$k_{s,y} \sim k_{i,y}^h \sim \tau_p Q$$

Factorization in the Coplanar Limit

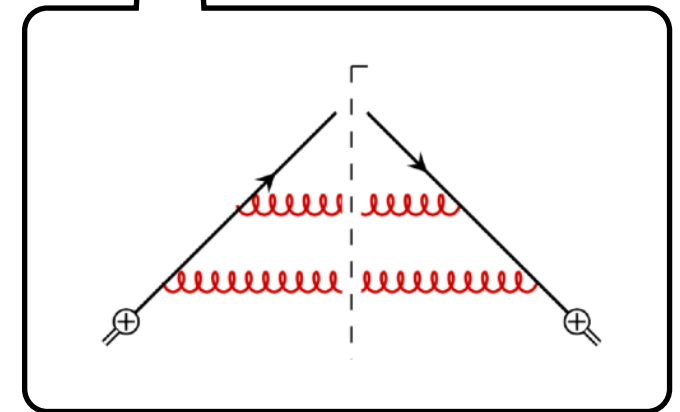
$$\frac{1}{\hat{\sigma}_0} \frac{d\sigma}{d\tau_p} = \int_D dv d\omega H(v, \omega, \mu) \int \frac{db}{2\pi\xi} 2 \cos(b\tau_p/\xi) S(b, \mu, \nu) J_q(b, \mu, \nu) J_{\bar{q}}(b, \mu, \nu) J_g(b, \mu, \nu)$$



Matrix Element of
3 jet production



VEV of semi-infinity
Wilson Loop



Second moment of quark/gluon
TMD fragmentation function

Probe the TMD gluon fragmentation at CEPC!

Resummation by RGE

RG

$$\frac{dH}{d \ln \mu^2} = \left[\frac{C_A + 2C_F}{2} \gamma_{\text{cusp}}(\alpha_s) \ln \frac{Q^2}{\mu^2} + \gamma_H(y, z, \alpha_s) \right] H ,$$

$$\begin{aligned} \frac{d \ln S}{d \ln \mu^2} = & \left[\frac{2C_F + C_A}{2} \left(\gamma_{\text{cusp}}[\alpha_s] \ln \frac{\mu^2}{\nu^2} - \gamma_s[\alpha_s] \right) + \frac{C_A}{2} \gamma_{\text{cusp}}[\alpha_s] \ln \frac{(1-u)^2 u}{v\omega} \right. \\ & \left. + C_F \gamma_{\text{cusp}}[\alpha_s] \ln \frac{(1-v)(1-\omega)}{u} \right] , \end{aligned}$$

$$\frac{d \ln S}{d \ln \nu^2} = \frac{2C_F + C_A}{2} \left(\int_{\mu^2}^{b_0^2/b^2} \frac{d\bar{\mu}^2}{\bar{\mu}^2} \gamma_{\text{cusp}}(\alpha_s[\bar{\mu}]) + \gamma_r(\alpha_s[b_0/b]) \right) ,$$

rapidity

RG

$$\frac{dJ_i}{d \ln \mu^2} = \left(-\frac{1}{2} C_i \gamma_{\text{cusp}} \ln \frac{(2p_i^0)^2}{\nu^2} + \gamma_{J,i} \right) J_i ,$$

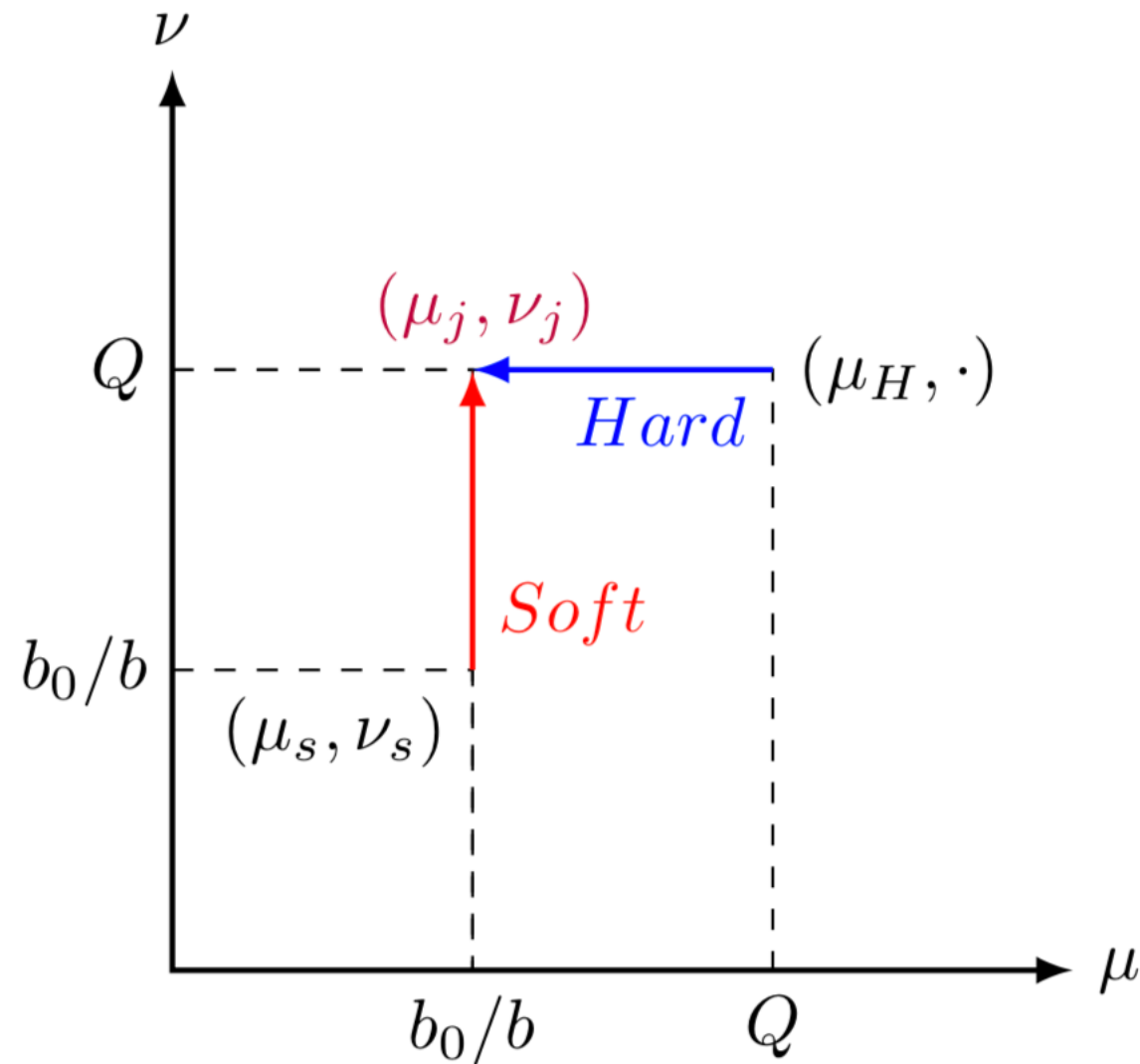
$$\frac{dJ_i}{d \ln \nu^2} = \frac{C_i}{2} \left(\int_{b_0^2/b^2}^{\mu^2} \frac{d\bar{\mu}^2}{\bar{\mu}^2} \gamma_{\text{cusp}}[\alpha_s(\bar{\mu})] - \gamma_r[\alpha_s(b_0/b)] \right) J_i .$$

- All the anomalous dimensions are known to at least three loops
- **RG invariant condition**

$$\gamma_H - \frac{C_A + 2C_F}{2} \gamma_s - 2\gamma_{J,q} - \gamma_{J,g} = 0$$

RG Evolution in both Virtuality and Rapidity

Setting $\mu = \mu_j = b_0/b$, $\nu = \nu_j = Q$,

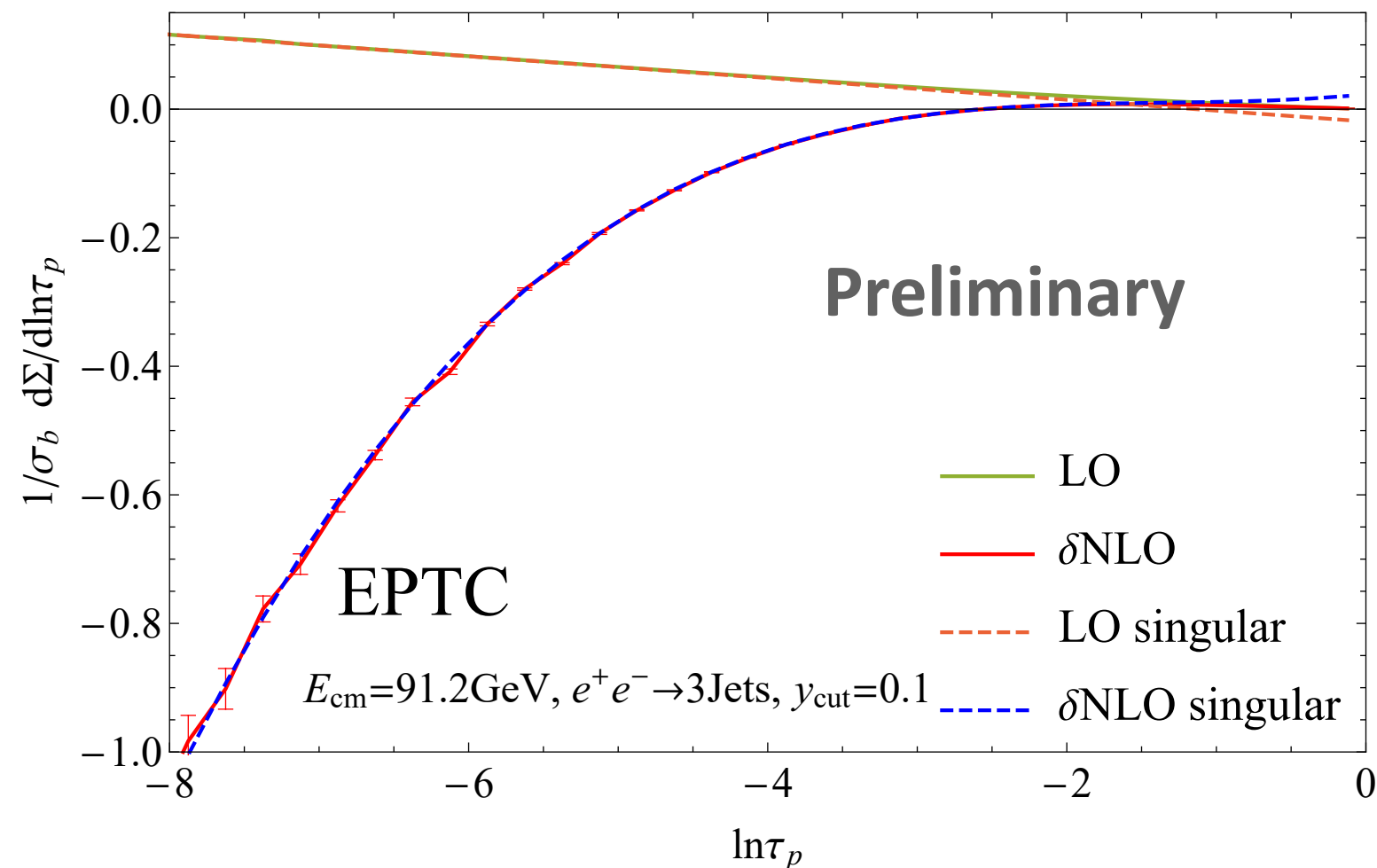


There is rapidity evolution for the soft function

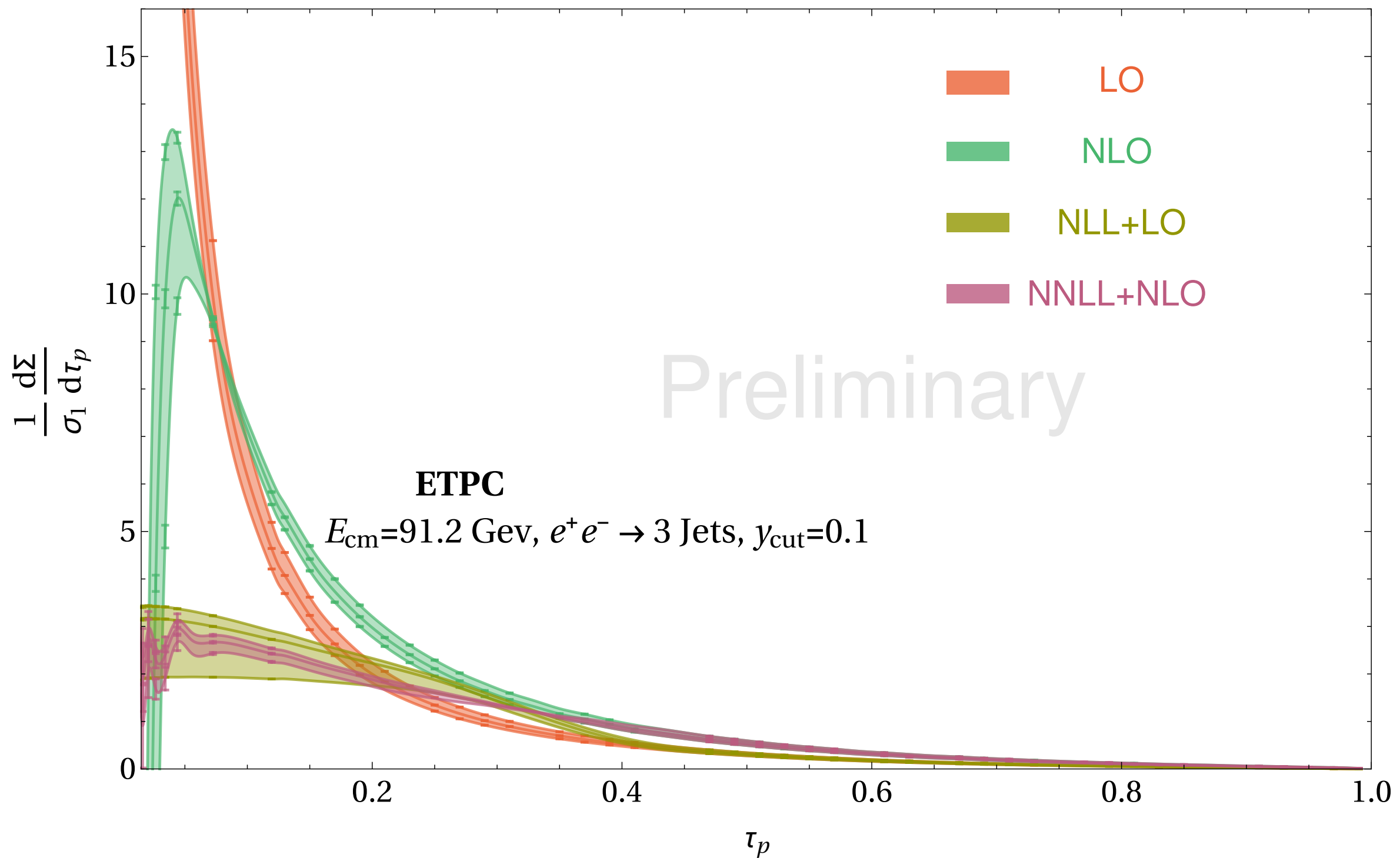
Validation of the Factorization Formula

$$\frac{d\Sigma}{d \ln \tau_p} = \int_D dv d\omega \left\{ \left(\frac{\alpha_s}{4\pi} \right)^2 (c_1 \ln \tau_p + c_2) + \left(\frac{\alpha_s}{4\pi} \right)^3 (c_3 \ln^3 \tau_p + c_4 \ln^2 \tau_p + c_5 \ln \tau_p + c_6) + \mathcal{O}(\tau_p^2) \right\}.$$

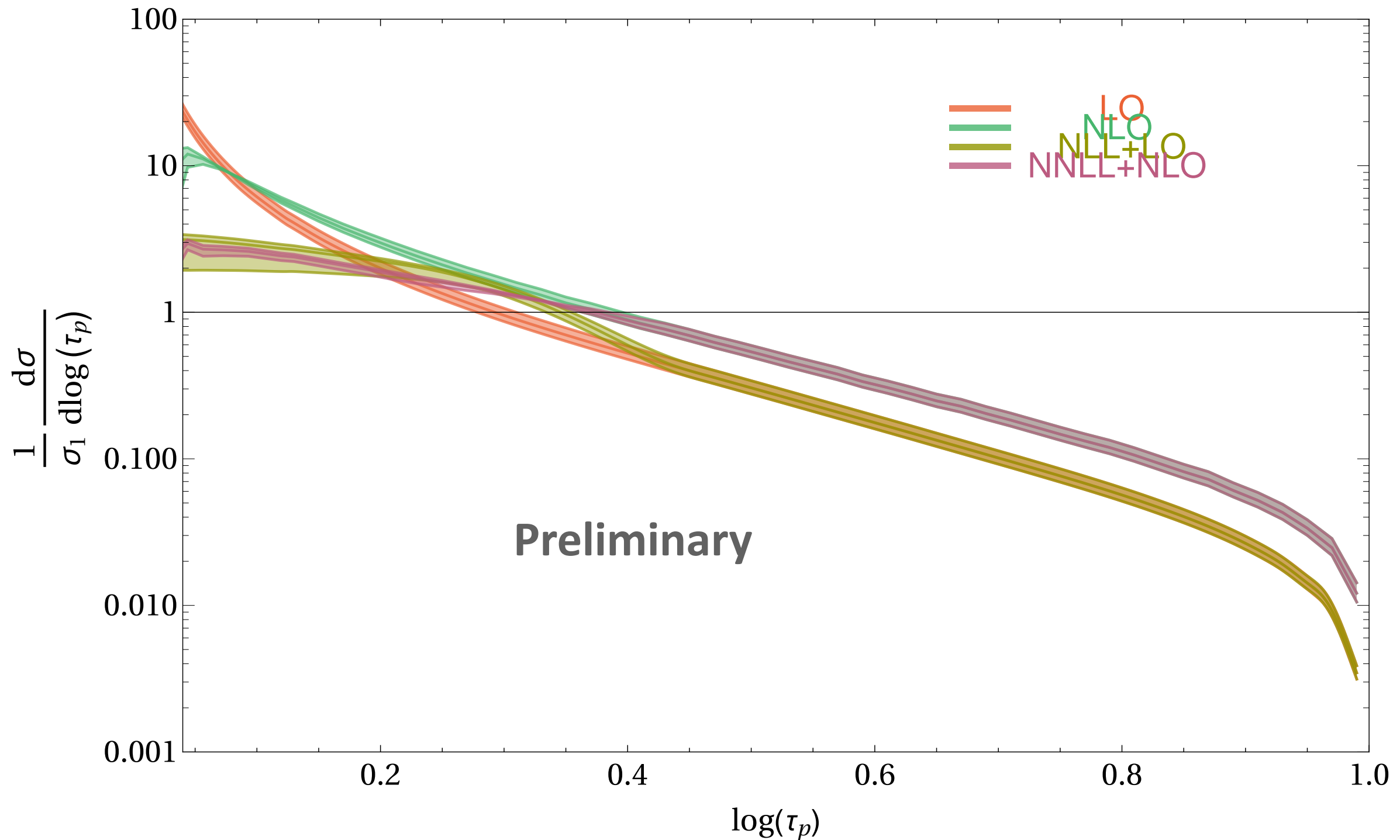
- Expand the factorization formula to NNLO, and compare with e+e- to 4 jet production at NLO (NLOJet++)



Resummation Prediction



Resummation Prediction



Summary

- A new three jet event shape: Energy Triple-Product Correlator
- Singular limit controlled by off-plane soft/collinear radiation
- All order factorization formula in terms well definite operator matrix element
- Application at the CEPC: α_s measurement, TMD fragmentation function, testing parton shower